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On some expectation and derivative operators related to integral representations of random variables with respect to a PII process.

STÉPHANE GOUTTE ^{*†}, NADIA OUDJANE [‡] AND FRANCESCO RUSSO [§]

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Abstract

Given a process with independent increments X (not necessarily a martingale) and a large class of square integrable r.v. $H = f(X_T)$, f being the Fourier transform of a finite measure μ , we provide explicit Kunita-Watanabe and Föllmer-Schweizer decompositions. The representation is expressed by means of two significant maps: the expectation and derivative operators related to the characteristics of X . We also provide an explicit expression for the variance optimal error when hedging the claim H with underlying process X . Those questions are motivated by finding the solution of the celebrated problem of global and local quadratic risk minimization in mathematical finance.

Key words and phrases: Föllmer-Schweizer decomposition, Kunita-Watanabe decomposition, Lévy processes, Characteristic functions, Processes with independent increments, global and local quadratic risk minimization, expectation and derivative operators.

2010 AMS-classification: 60G51, 60H05, 60J75, 91G10

1 Introduction

Let X be an (\mathcal{F}_t) -special cadlag semimartingale, where (\mathcal{F}_t) is a filtration fulfilling the usual conditions. It admits a unique decomposition $M + A$ where M is an (\mathcal{F}_t) -local martingale and A is an (\mathcal{F}_t) -predictable process with bounded variation. Given $T > 0$ and a square integrable random variable H which is \mathcal{F}_T -measurable, we consider three specific issues of stochastic analysis that are particularly relevant in stochastic finance.

Kunita-Watanabe (KW) decomposition. This problem consists in providing existence conditions and explicit expressions of a predictable process $(Z_t)_{t \in [0, T]}$ and an \mathcal{F}_0 -measurable r.v. such that

$$H = V_0 + \int_0^T Z_s dM_s + O_T, \quad (1.1)$$

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where (O_t) is an (\mathcal{F}_t) -local martingale such that $\langle O, M \rangle = 0$.

When $X = M$ is a classical Brownian motion W and (\mathcal{F}_t) is the associated canonical filtration, Z is provided by the celebrated *Clark-Ocone formula* at least when H belongs to the Malliavin-Sobolev type space $\mathbb{D}^{1,2}$. In that case one has

$$H = \mathbb{E}(H) + \int_0^T \mathbb{E}(D_s H | \mathcal{F}_s) dW_s, \quad (1.2)$$

where $DH = (D_t H)_{t \in [0, T]}$ is the classical Malliavin derivative of H .

In the last ten years a significant scientific production appeared at the level of Malliavin calculus in relation with Poisson measures in several directions. A trend which was particularly directed to obtaining a generalization of Clark-Ocone formula was started by [22]. In Theorem 1, the authors obtained a chaos type decomposition of a square integrable random variable H in the Poisson space generated by a finite number of Lévy square integrable martingales (η_j) , with respect to a well-chosen sequence of strongly orthogonal martingales $\gamma^{(m)}$. This could allow to represent any H as an infinite sum of stochastic integrals with respect to the $\gamma^{(m)}$, an infinite dimensional *derivative* $\mathcal{D}^{(m)}$ with respect to $\gamma^{(m)}$ and a Malliavin-Sobolev type space $\mathbb{D}^{1,2}$. A first formulation of a Clark-Ocone type formula was produced by [19]: it consisted in representing square integrable random variables H with respect to the $\gamma^{(m)}$ in terms of some predictable projections of $\mathcal{D}^{(m)} H$. Another class of stochastic derivative (this time) with respect to η_j was introduced by [9]. With the help of an isometry obtained in [20], one could deduce the more intrinsic (and recently widely used) Clark-Ocone type formula of the type

$$H = \mathbb{E}(H) + \int_0^T \int_{\mathbb{R}} \mathbb{E}(D_{t,x} H | \mathcal{F}_t) \tilde{N}(dt, dx)$$

where \tilde{N} is the compensated Poisson random measure and $(D_{t,x})$ a two-indexed derivative operator. This formula is also stated in Theorem 12.16 of [10]. Theorem 4.1 of [3] allows to provide an explicit representation of the process Z appearing in (1.1) with the help of previous operator $D_{t,x}$.

Föllmer-Schweizer decomposition. That decomposition is a generalization of the Kunita-Watanabe one in the sense that square integrable random variables are represented with respect to X instead of M . It consists in providing existence conditions and explicit expressions of a predictable process ξ and an \mathcal{F}_0 -measurable square integrable r.v. H_0 such that

$$H = H_0 + \int_0^T \xi_s dX_s + L_T \quad (1.3)$$

where L_T is the terminal value of an orthogonal martingale L to M , the martingale part of X .

In the seminal paper [12], the problem is treated for an underlying process X with continuous paths. In the general case, X is said to satisfy the **structure condition** (SC) if there is a predictable process α such that $A_t = \int_0^t \alpha_s d\langle M \rangle_s$, $t \in [0, T]$, and $\int_0^T \alpha_s^2 d\langle M \rangle_s < \infty$ a.s. An interesting connection with the theory of backward stochastic differential equations (BSDEs) in the sense of [23], was proposed in [27]. [23] considered BSDEs driven by Brownian motion; in [27] the Brownian motion is in fact replaced by M . The first author who considered a BSDE driven by a martingale was [5]. The BSDE problem consists in finding a triple (V, ξ, L) where

$$V_t = H - \int_t^T \xi_s dM_s - \int_t^T \xi_s \alpha_s d\langle M \rangle_s - (L_T - L_t),$$

and L is an (\mathcal{F}_t) -local martingale orthogonal to M . The solution (V_0, ξ, L) of that BSDE constitutes a triplet (H_0, ξ, L) solving (1.3). The FS decomposition is motivated in mathematical finance by looking for the solution of the so called *local risk minimization*, see [12] where H represents a contingent claim to hedge and X is related to the price of the underlying asset. In this case, V_t represents the *hedging portfolio value* of the contingent claim at time t , ξ represents the hedging strategy and the initial capital V_0 constitutes in fact the expectation of H under the so called *minimal martingale measure*, see [28].

Variance optimal hedging. This approach developed by M. Schweizer ([27], [29]) consists in minimizing the quadratic distance between the hedging portfolio and the pay-off. More precisely, it consists in providing existence conditions and explicit expressions of a predictable process $(\varphi_t)_{t \in [0, T]}$ and an \mathcal{F}_0 -measurable square integrable r.v. V_0 such that

$$(V_0, \varphi) = \text{Argmin}_{c, v} \mathbb{E}(\varepsilon(c, v))^2, \quad \text{where} \quad \varepsilon(c, v) = H - c - \int_0^T v_s dX_s. \quad (1.4)$$

The quantity V_0 and process φ represent the initial capital and the optimal hedging strategy of the contingent claim H .

When the market is complete and without arbitrage opportunities, the representation property (1.3) holds with $L \equiv 0$; so those three decompositions (Kunita-Watanabe, Föllmer-Schweizer and Variance Optimal) reduce to a single representation of the random variable H as a stochastic integral modulo a martingale (*risk neutral*) change of measure. If the market model is incomplete (e.g. because of jumps or stochastic volatility in prices dynamics) then those three decompositions are in general different and a residual term must be added to each integral representation, e.g. O_T and L_T and $\varepsilon(V_0, \varphi)$. However, even in this incomplete market setting, a nice exception occurs if the underlying price X is a martingale. Indeed, the martingale property allows to bypass some theoretical difficulties leading again to three identical decompositions.

Most of the articles providing quasi-explicit expressions for those decompositions are precisely assuming the martingale property for the process X , therefore coming down to consider the Kunita-Watanabe decomposition. For instance in [16], the authors developed an original approach to find an explicit expression for the Kunita-Watanabe decomposition of a random variable H of the form $H = f(Y_T)$ where Y is a reference Markov process and the price process X is a martingale related to Y . Their idea is to apply Ito's formula to derive the Doob-Meyer decomposition of $\mathbb{E}[H|\mathcal{F}_t]$ and then to write the orthogonality condition between $\mathbb{E}[H|\mathcal{F}_t] - \int_0^t Z_s dX_s$ and X . In [7], the authors follow the same idea to derive the hedging strategy minimizing the *Variance Optimal hedging error* under the (risk-neutral) pricing measure. They provide some interesting financial motivations for this martingale framework. Their approach also applies to a broad class of price models and to some path dependent random variables H . In some specific cases they obtain quasi-explicit expressions for the Variance Optimal strategy. For instance, they prove that if X is the exponential of a Lévy process, then the strategy is related to derivatives and integrals w.r.t. the Lévy measure of the conditional expectation $\mathbb{E}[H|\mathcal{F}_t]$.

Unfortunately, minimizing the quadratic hedging error under the pricing measure, can lead to a huge quadratic error under the objective measure. Moreover, the use of Ito's lemma in those approaches requires some regularity conditions on the conditional expectation $\mathbb{E}[H|\mathcal{F}_t]$: basically it should be once differentiable w.r.t. the time variable and twice differentiable w.r.t. the space variable with continuous partial derivatives.

In the non-martingale framework, one major contribution is due to [15] whose authors restricted their analysis to the specific case where X is the exponential of a Lévy process and $H = f(X_T)$, f being the

Fourier-Laplace transform of a complex finite measure. The authors obtained an explicit expression for the process ξ intervening in (1.3). This result was generalized to exponential of non stationary processes in the continuous and discrete time setting in [13] and [14].

Following this approach, the objective of the present paper is to consider the non-martingale framework and to provide quasi explicit expressions of both the Kunita-Watanabe and Föllmer-Schweizer decompositions when X is a general process with independent increments and $H = f(X_T)$ is the Fourier transform of a finite measure μ . Our method does not rely on Ito's formula and therefore does not require any further regularity condition on conditional expectations. The representation is carried by means of two significant maps: the so-called *expectation and derivative operators* related to the characteristics of the underlying process X . We also express explicitly the Variance Optimal hedging strategy and the corresponding Variance Optimal error.

The paper is organized as follows. In Section 2 we recall some essential considerations related to the Föllmer-Schweizer decomposition related to general special semimartingales. In Section 3 we provide the framework related to processes with independent increments and related structure conditions. Section 4 provides the explicit Kunita-Watanabe and the Föllmer-Schweizer decompositions under minimal assumptions. Section 5 formulates the solution of the global minimization problem evaluating the variance of the hedging error. Finally, in Section 6, we consider a class of examples, for which we verify that the assumptions are fulfilled.

2 Generalities on Föllmer-Schweizer decomposition and mean variance hedging

In the whole paper, $T > 0$, will be a fixed terminal time and we will denote by $(\Omega, \mathcal{F}, (\mathcal{F}_t)_{t \in [0, T]}, P)$ a filtered probability space, fulfilling the usual conditions. We suppose from now on \mathcal{F}_0 to be trivial for simplicity.

2.1 Optimality and Föllmer-Schweizer Structure Condition

Let $X = (X_t)_{t \in [0, T]}$ be a real-valued special semimartingale with canonical decomposition, $X = M + A$. For the clarity of the reader, we formulate in dimension one, the concepts appearing in the literature, see e.g. [27] in the multidimensional case. In the sequel Θ will denote the space $L^2(M)$ of all predictable \mathbb{R} -valued processes $v = (v_t)_{t \in [0, T]}$ such that $\mathbb{E} \left[\int_0^T |v_s|^2 d \langle M \rangle_s \right] < \infty$. For such v , clearly $\int_0^t v dX, t \in [0, T]$ is well-defined; we denote by $G_T(\Theta)$, the space generated by all the r.v. $G_T(v) = \int_0^T v_s dX_s$ with $v = (v_t)_{t \in [0, T]} \in \Theta$.

Definition 2.1. *The **minimization problem** we aim to study is the following: Given $H \in \mathcal{L}^2$, an admissible strategy pair (V_0, φ) will be called **optimal** if $(c, v) = (V_0, \varphi)$ minimizes the expected squared hedging error*

$$\mathbb{E}[(H - c - G_T(v))^2], \quad (2.1)$$

*over all admissible strategy pairs $(c, v) \in \mathbb{R} \times \Theta$. V_0 will represent the **initial capital** of the hedging portfolio for the contingent claim H at time zero.*

The definition below introduces an important technical condition, see [27].

Definition 2.2. *Let $X = (X_t)_{t \in [0, T]}$ be a real-valued special semimartingale. X is said to satisfy the **structure condition (SC)** if there is a predictable \mathbb{R} -valued process $\alpha = (\alpha_t)_{t \in [0, T]}$ such that the following properties are verified.*

1. $A_t = \int_0^t \alpha_s d\langle M \rangle_s$, for all $t \in [0, T]$, so that $dA \ll d\langle M \rangle$.
2. $\int_0^T \alpha_s^2 d\langle M \rangle_s < \infty$, P -a.s.

Definition 2.3. From now on, we will denote by $K = (K_t)_{t \in [0, T]}$ the cadlag process $K_t = \int_0^t \alpha_s^2 d\langle M \rangle_s$, for all $t \in [0, T]$. This process will be called the **mean-variance trade-off (MVT)** process.

In [27], the process $(K_t)_{t \in [0, T]}$ is denoted by $(\hat{K}_t)_{t \in [0, T]}$.

2.2 Föllmer-Schweizer decomposition and variance optimal hedging

Throughout this section, as in Section 2.1, X is supposed to be an (\mathcal{F}_t) -special semimartingale fulfilling the (SC) condition.

Definition 2.4. We say that a random variable $H \in \mathcal{L}^2(\Omega, \mathcal{F}, P)$ admits a **Föllmer-Schweizer (FS) decomposition**, if it can be written as

$$H = H_0 + \int_0^T \xi_s^H dX_s + L_T^H, \quad P - a.s., \quad (2.2)$$

where $H_0 \in \mathbb{R}$ is a constant, $\xi^H \in \Theta$ and $L^H = (L_t^H)_{t \in [0, T]}$ is a square integrable martingale, with $\mathbb{E}[L_0^H] = 0$ and strongly orthogonal to M , i.e. $\langle L^H, M \rangle = 0$.

The notion of strong orthogonality is treated for instance in Chapter IV.3 p. 179 of [24].

Theorem 2.5. If X satisfies (SC) and the MVT process K is uniformly bounded in t and ω , then we have the following.

1. Every random variable $H \in \mathcal{L}^2(\Omega, \mathcal{F}, P)$ admits a unique FS decomposition. Moreover, $H_0 \in \mathbb{R}$, $\xi \in \Theta$ and L^H is uniquely determined by H .
2. For every $H \in \mathcal{L}^2(\Omega, \mathcal{F}, P)$ there exists a unique $(c^{(H)}, \varphi^{(H)}) \in \mathbb{R} \times \Theta$ such that

$$\mathbb{E}[(H - c^{(H)} - G_T(\varphi^{(H)}))^2] = \min_{(c, v) \in \mathbb{R} \times \Theta} \mathbb{E}[(H - c - G_T(v))^2]. \quad (2.3)$$

From the Föllmer-Schweizer decomposition follows the solution to the global minimization problem (2.1). Next theorem gives the explicit form of the optimal strategy.

Theorem 2.6. Suppose that X satisfies (SC), that the MVT process K of X is deterministic and $\langle M \rangle$ is continuous. Let α be the process appearing in Definition 2.2 of (SC) and let $H \in \mathcal{L}^2$.

$$\min_{(c, v) \in \mathbb{R} \times \Theta} \mathbb{E}[(H - c - G_T(v))^2] = \exp(-K_T) \mathbb{E}[(L_0^H)^2] + \mathbb{E} \left[\int_0^T \exp\{-(K_T - K_s)\} d\langle L^H \rangle_s \right].$$

Proof. The result follows from Corollary 9 of [27]. We remark that being $\langle M \rangle$ continuous, the Doléans-Dade exponential of K , $\mathcal{E}(K)$, equals $\exp(K)$. \square

In the sequel, we will find an explicit expression of the KW and FS decomposition for a large class of square integrable random variables H , when the underlying process is a process with independent increments.

3 Processes with independent increments (PII)

This section deals with the case of processes with independent increments. First, we recall some useful properties of such processes, then, we obtain a sufficient condition on the characteristic function for the existence of the FS decomposition.

Beyond its own theoretical interest, this work is motivated by its possible application to hedging derivatives related to financial or commodity assets. Indeed, in some specific cases it is reasonable to introduce arithmetic models (eg. Bachelier) in contrast to geometric models (eg. Black-Scholes model), see for instance [2].

3.1 Generalities on PII processes

Let $X = (X_t)_{t \in [0, T]}$ be a stochastic process. Let $t \in [0, T]$.

Definition 3.1. 1. The **characteristic function** of (the law of) X_t is the continuous function

$$\varphi_t : \mathbb{R} \rightarrow \mathbb{C} \quad \text{with} \quad \varphi_t(u) = \mathbb{E}[e^{iuX_t}].$$

2. The **Log-characteristic function** of (the law of) X_t is the unique function $\Psi_t : \mathbb{R} \rightarrow \mathbb{R}$ such that $\varphi_t = \exp(\Psi_t(u))$ and $\Psi_t(0) = 0$.

Notice that for $u \in \mathbb{R}$ we have $\overline{\Psi_t(u)} = \Psi_t(-u)$. Since $\varphi : [0, T] \times \mathbb{R} \rightarrow \mathbb{C}$, is uniformly continuous and $\varphi_t(0) = 1$, then there is a neighborhood \mathcal{U} of 0 such that

$$\operatorname{Re} \Psi_t(u) > 0, \forall t \in [0, T], u \in \mathcal{U}. \quad (3.1)$$

Definition 3.2. $X = (X_t)_{t \in [0, T]}$ is a (real) **process with independent increments (PII)** iff

1. X has cadlag paths;
2. $X_0 = 0$;
3. $X_t - X_s$ is independent of \mathcal{F}_s for $0 \leq s < t \leq T$ where (\mathcal{F}_t) is the canonical filtration associated with X ; moreover we will also suppose
4. X is continuous in probability, i.e. X has no fixed time of discontinuities.

The process X is said to be **square integrable** if for every $t \in [0, T]$, $\mathbb{E}[|X_t|^2] < \infty$.

From now on, (\mathcal{F}_t) will always be the canonical filtration associated with X . Below, we state some elementary properties of the characteristic functions related to PII processes. In the sequel, we will always suppose that X is a semimartingale. For more details about those processes the reader can consult Chapter II of [17].

Remark 3.3. Let $0 \leq s < t \leq T$, $u \in \mathbb{R}$,

1. $\Psi_{X_t}(u) = \Psi_{X_s}(u) + \Psi_{X_t - X_s}(u)$.
2. $\exp(iuX_t - \Psi_t(u))$ is an (\mathcal{F}_t) -martingale.

3. There is an increasing function $a : [0, T] \rightarrow \mathbb{R}$ and a triplet (b_t, c_t, F_t) called characteristics such that

$$\Psi_t(u) = \int_0^t \eta_s(u) da_s, \quad \text{for all } u \in \mathbb{R}. \quad (3.2)$$

where $\eta_s(u) := \left[i u b_s - \frac{u^2}{2} c_s + \int_{\mathbb{R}} (e^{i u x} - 1 - i u x \mathbf{1}_{|x| \leq 1}) F_s(dx) \right]$. Indeed $b : [0, T] \rightarrow \mathbb{R}$, $c : [0, T] \rightarrow \mathbb{R}_+$ are deterministic functions and for any $t \in [0, T]$, F_t is a positive measure such that

$\int_{[0, T] \times \mathbb{R}} (1 \wedge x^2) F_t(dx) da_t < \infty$. For more details we refer to the statement and the proof of Proposition II.2.9 of [17].

4. The Borel measure on $[0, T] \times \mathbb{R}$ defined by $F_t(dx) da_t$ is called jump measure and it is denoted by $\nu(dt, dx)$.

5. We have $\int_{[0, T] \times \mathbb{R}} x^2 \nu(dt, dx) = \mathbb{E} \left(\sum_{t \in [0, T]} (\Delta X_t)^2 \right)$ where $\Delta X_t = X_t - X_{t-}$ is the jump at time t of the process X .

6. Suppose that X is square integrable. Since previous sum of jumps is bounded by the square bracket at time T , i.e. $[X, X]_T$, which is integrable, it follows that $\mathbb{E} \left(\sum_{t \in [0, T]} (\Delta X_t)^2 \right) < \infty$.

Remark 3.4. 1. The process X is square integrable if and only if for every $t \in [0, T]$, $u \mapsto \varphi_t(u)$ is of class C^2 .

2. By (3.1), X is square integrable if and only if $u \mapsto \Psi_t(u)$ is of class C^2 , $t \in [0, T]$, $u \in \mathcal{U}$.

3. If X is square integrable, the chain rule derivation implies

$$\mathbb{E}[X_t] = -i \Psi'_t(0), \quad \mathbb{E}[X_t - X_s] = -i(\Psi'_t(0) - \Psi'_s(0)), \quad (3.3)$$

$$\text{Var}(X_t) = -\Psi''_t(0), \quad (3.4)$$

$$\text{Var}(X_t - X_s) = -[\Psi''_t(0) - \Psi''_s(0)]. \quad (3.5)$$

Remark 3.5. Suppose that X is a square integrable PII process. We observe that it is possible to permute integral and derivative in the expression (3.2). In fact consider $t \in [0, T]$. We need to show that

$$\frac{d}{du} \int_{[0, t] \times \mathbb{R}} \nu(ds, dx) g(s, x; u) = \int_{[0, t] \times \mathbb{R}} \nu(ds, dx) \frac{\partial g}{\partial u}(s, x; u), \quad (3.6)$$

where $g(s, x; u) = ix(e^{iux} - \mathbf{1}_{|x| \leq 1})$. We observe that

$$\begin{aligned} |g(s, x; u)|^2 &= |e^{iux} - \mathbf{1}_{|x| \leq 1}|^2 = |\cos ux - \mathbf{1}_{|x| \leq 1}|^2 + |\sin ux|^2 = \mathbf{1}_{|x| > 1} + 4\mathbf{1}_{|x| \leq 1} \left(\sin \frac{ux}{2}\right)^2 \\ &\leq 4\left(\frac{u^2}{2} \vee 1\right)(x^2 \wedge 1). \end{aligned}$$

Hence, we obtain that for any real interval $[a, b]$ and for any $u \in [a, b]$,

$$\begin{aligned} \left| \frac{\partial g(s, x; u)}{\partial u} \right| &= \left| \frac{\partial}{\partial u} (e^{iux} - 1 - iux \mathbf{1}_{|x| \leq 1}) \right| = |ix(e^{iux} - \mathbf{1}_{|x| \leq 1})| \leq 2x^2(|u| \vee 1) \\ &\leq 2bx^2 =: \gamma(s, x). \end{aligned}$$

Consequently by finite increments theorem, Remark 3.3 5) it follows that $\int_{[0, T] \times \mathbb{R}} \nu(ds, dx) \gamma(s, x) < \infty$. By the definition of derivative and Lebesgue dominated convergence theorem the result (3.6) follows. So

$$\Psi'_t(u) = i \int_0^t b_s da_s - u \int_0^t c_s da_s + \int_0^t \left(\int_{\mathbb{R}} ix(e^{iux} - \mathbf{1}_{|x| \leq 1}) F_s(dx) \right) da_s, \quad \text{for all } u \in \mathbb{R}. \quad (3.7)$$

Moreover, since

$$\left| \frac{\partial}{\partial u} (ix(e^{iux} - \mathbf{1}_{|x| \leq 1})) \right| = |x^2 e^{iux}| \leq x^2$$

we obtain similarly

$$\Psi_t''(u) = - \int_0^t c_s da_s - \int_{[0,T] \times \mathbb{R}} x^2 e^{iux} F_s(dx) da_s = - \int_0^t \xi_s(u) da_s, \quad (3.8)$$

where $\xi_s(u) = c_s + \int_{\mathbb{R}} x^2 e^{iux} F_s(dx)$. In particular, for every $u \in \mathbb{R}$, $t \mapsto \Psi_t'(u)$ and $t \mapsto \Psi_t''(u)$ are absolutely continuous with respect to da_s .

Remark 3.6. Suppose that X is square integrable. A consequence of Remark 3.5 is the following.

1. $t \mapsto \Psi_t'(u)$ is continuous for every $u \in \mathbb{R}$ and therefore bounded on $[0, T]$.
2. $t \mapsto \Psi_t''(0)$ is continuous.

3.2 Structure condition for PII

Let $X = (X_t)_{t \in [0, T]}$ be a real-valued semimartingale with independent increments and $X_0 = 0$. From now on, X will be supposed to be **square integrable**.

Proposition 3.7. 1. X is a special semimartingale with decomposition $X = M + A$ with the following properties:

$\langle M \rangle_t = -\Psi_t''(0)$ and $A_t = -i\Psi_t'(0)$. In particular $t \mapsto -\Psi_t''(0)$ is increasing and therefore of bounded variation.

2. X satisfies condition (SC) of Definition 2.2 if and only if

$$d\Psi_t'(0) \ll d\Psi_t''(0) \quad \text{and} \quad \int_0^T \left| \frac{d\Psi_s'(0)}{d\Psi_s''(0)} \right|^2 |d\Psi_s''(0)| < \infty. \quad (3.9)$$

In that case

$$A_t = \int_0^t \alpha_s d\langle M \rangle_s \quad \text{with} \quad \alpha_t = i \frac{d\Psi_t'(0)}{d\Psi_t''(0)} \quad \text{for all } t \in [0, T]. \quad (3.10)$$

3. Under condition (3.9), FS decomposition exists (and it is unique) for every square integrable random variable.

Before going into the proof of the above proposition, let us derive one implication on the validity of the (SC) in the Lévy case. Let $X = (X_t)_{t \in [0, T]}$ be a real-valued Lévy process, with $X_0 = 0$. We assume that $\mathbb{E}[|X_T|^2] < \infty$.

1. Since $X = (X_t)_{t \in [0, T]}$ is a Lévy process then $\Psi_t(u) = t\Psi_1(u)$. In the sequel, we will use the shortened notation $\Psi := \Psi_1$.
2. Ψ is a function of class C^2 and $\Psi''(0) = \text{Var}(X_1)$ which is strictly positive if X_1 is non deterministic.

Then by application of Proposition 3.7, we get the following result.

Corollary 3.8. Let $X = M + A$ be the canonical decomposition of the Lévy process X . Then for all $t \in [0, T]$,

$$\langle M \rangle_t = -t\Psi''(0) \quad \text{and} \quad A_t = -it\Psi'(0). \quad (3.11)$$

If $\Psi''(0) \neq 0$ then X satisfies condition (SC) of Definition 2.2 with

$$A_t = \int_0^t \alpha d\langle M \rangle_s \quad \text{with} \quad \alpha = i \frac{\Psi'(0)}{\Psi''(0)} \quad \text{for all } t \in [0, T]. \quad (3.12)$$

Hence, FS decomposition exists for every square integrable random variable. If $\Psi''(0) = 0$ then (X_t) verifies condition (SC) if and only if $X_t \equiv 0$.

Proof of Proposition 3.7. 1. Let us first determine A and M in terms of the log-characteristic function of X . Using (3.3) of Remark 3.4, we get $\mathbb{E}[X_t | \mathcal{F}_s] = \mathbb{E}[X_t - X_s + X_s | \mathcal{F}_s] = -i\Psi'_t(0) + i\Psi'_s(0) + X_s$, then $\mathbb{E}[X_t + i\Psi'_t(0) | \mathcal{F}_s] = X_s + i\Psi'_s(0)$. Hence, $(X_t + i\Psi'_t(0))$ is a martingale and the canonical decomposition of X follows $X_t = \underbrace{X_t + i\Psi'_t(0)}_{M_t} - \underbrace{i\Psi'_t(0)}_{A_t}$, where M is a local martingale and A is a locally bounded variation process thanks to the semimartingale property of X . Let us now determine $\langle M \rangle$, in terms of the log-characteristic function of X . Using (3.3) and (3.5) of Remark 3.4, yields

$$\begin{aligned} \mathbb{E}[M_t^2 | \mathcal{F}_s] &= \mathbb{E}[(X_t + i\Psi'_t(0))^2 | \mathcal{F}_s] = \mathbb{E}[(M_s + X_t - X_s + i(\Psi'_t(0) - \Psi'_s(0)))^2 | \mathcal{F}_s], \\ &= M_s^2 + \text{Var}(X_t - X_s) = M_s^2 - \Psi''_t(0) + \Psi''_s(0). \end{aligned}$$

Hence, $(M_t^2 + \Psi''_t(0))$ is a (\mathcal{F}_t) -martingale, and point 1. is established.

2. is a consequence of point 1. and of Definition 2.2. On the other hand $A_t = \int_0^t \alpha_s d\langle M \rangle_s$ with $\alpha_t = i \frac{d\Psi'_t(0)}{d\Psi''_t(0)}$ for $t \in [0, T]$.

3. follows from Theorem 2.5. In fact $K_T = \int_0^T \left(\frac{d\Psi'_s(0)}{d\Psi''_s(0)} \right)^2 d(-\Psi''_s(0))$ is deterministic and in particular K is uniformly bounded. □

Condition (SC) implies a significant necessary condition.

Proposition 3.9. *If X satisfies condition (SC), then one of the two following properties hold.*

1. X has no deterministic increments.
2. If $X_b - X_a$ is deterministic then $X_u = X_a, \forall u \in [a, b]$.

Proof. We suppose that (SC) is fulfilled and let $0 \leq a < b \leq T$ for which $X_b - X_a$ is deterministic. Consequently $-(\Psi''_b(0) - \Psi''_a(0)) = \text{Var}(X_b - X_a) = 0$. This implies that $X_t - X_a$ is deterministic for every $t \in [a, b]$. By (3.9), it follows that $\Psi'_t(0) = \Psi'_a(0), \forall t \in [a, b]$. Hence, for any $t \in [a, b]$, we have $X_t - X_a = \mathbb{E}[X_t - X_a] = 0$. □

The following technical result will be useful in the sequel.

Proposition 3.10. *If X satisfies condition (SC), there is $\tilde{a} : [0, T] \rightarrow \mathbb{R}$ increasing such that $d\tilde{a}_t$ is equivalent to $-d(\Psi''_t(0))$ and (3.2) holds with da_t replaced by $d\tilde{a}_t$.*

Proof. Appendix. □

From now on, a_t will be replaced by $-\Psi''_t(0)$. Equalities and inequalities will generally hold $d(-\Psi''_t(0))$ a.e. with respect to t .

Corollary 3.11. *We suppose that X is square integrable and it fulfills (SC). Then for every $u \in \mathbb{R}$, $t \mapsto \Psi_t(u)$, $t \mapsto \Psi'_t(u)$ and $t \mapsto \Psi''_t(u)$, are a.c. w.r.t. $-\Psi''_t(0)$.*

1. In particular,

$$\Psi'_t(u) = \int_0^t \zeta_s(u) d(-\Psi''_s(0)) \quad \text{and} \quad \Psi''_t(u) = \int_0^t \xi_s(u) d(-\Psi''_s(0)),$$

where

$$\zeta_s(u) = ib_s - uc_s + \int_{\mathbb{R}} ix(e^{iux} - 1_{\{|x| \leq 1\}}) F_s(dx) \quad \text{and} \quad \xi_s(u) = c_s + \int_{\mathbb{R}} x^2 e^{iux} F_s(dx).$$

2. Setting $u = 0$, we obtain $\xi_s(0) = c_s + \int_{\mathbb{R}} x^2 F_s(dx) = 1$, $d(-\Psi''_s(0))$ a.e.

Proof. It follows from Proposition 3.10, item 3. of Remark 3.3 and Remark 3.5. □

3.3 Examples

3.3.1 A Gaussian continuous process example

Let $\psi : [0, T] \rightarrow \mathbb{R}$ be a continuous increasing function, $\gamma : [0, T] \rightarrow \mathbb{R}$ be a bounded variation function. We set $X_t = W_{\psi(t)} + \gamma(t)$, where W is the standard Brownian motion on \mathbb{R} . Clearly, $X_t = M_t + \gamma(t)$, where $M_t = W_{\psi(t)}$, defines a continuous martingale, such that $\langle M \rangle_t = [M]_t = \psi(t)$. Since $X_t \sim \mathcal{N}(\gamma(t), \psi(t))$, for all $u \in \mathbb{R}$ and $t \in [0, T]$, we have $\Psi_t(u) = i\gamma(t)u - \frac{u^2 \psi(t)}{2}$ which yields $\Psi'_t(0) = i\gamma(t)$ and $\Psi''_t(0) = -\psi(t)$. Taking into account Proposition 3.7 2, (SC) is verified if and only if $\gamma \ll \psi$ and $\frac{d\gamma}{d\psi} \in \mathcal{L}^2(d\psi)$. This is of course always verified if $\gamma \equiv 0$. We have $A_t = \int_0^t \alpha_s d\langle M \rangle_s$ and $\alpha_t = \frac{d\gamma}{d\psi} \Big|_t$ for all $t \in [0, T]$.

3.3.2 Processes with independent and stationary increments (Lévy processes)

We recall some log-characteristic functions of typical Lévy processes. In this case we have $\Psi_t(u) = t\Psi(u)$, $t \in [0, T]$, $x \in \mathbb{R}$.

1. **Poisson Case:** If X is a Poisson process with intensity λ , then for all $u \in \mathbb{R}$, $\Psi(u) = \lambda(e^{iu} - 1)$, $\Psi'(0) = i\lambda$ and $\Psi''(0) = -\lambda$, which yields $\alpha \equiv 1$.
2. **NIG Case:** This process was introduced by Barndorff-Nielsen in [1]. If X is a Normal Inverse Gaussian Lévy process with $X_1 \sim NIG(\theta, \beta, \delta, \mu)$, with $\theta > |\beta| > 0$, $\delta > 0$ then for all $u \in \mathbb{R}$, $\Psi(u) = \mu iu + \delta(\gamma_0 - \gamma_{iu})$, where $\gamma_{iu} = \sqrt{\theta^2 - (\beta + iu)^2}$. By derivation, one gets $\Psi'(0) = i\mu + \delta \frac{i\beta}{\gamma_0}$ and $\Psi''(0) = -\delta(\frac{1}{\gamma_0} + \frac{\beta^2}{\gamma_0^3})$ which yields $\alpha \equiv i \frac{\Psi'(0)}{\Psi''(0)} = \frac{\gamma_0^2(\gamma_0\mu + \delta\beta)}{\delta(\gamma_0^2 + \beta)}$.
3. **Variance Gamma case:** If X is a Variance Gamma process with $X_1 \sim VG(\theta, \beta, \delta, \mu)$ where $\theta, \beta > 0$, $\delta \neq 0$, then for all $u \in \mathbb{R}$, The expression of the log-characteristic function can be found in [15] or also [6], table IV.4.5 in the particular case $\mu = 0$. We have $\Psi(u) = \mu iu + \delta \text{Log} \left(\frac{\theta}{\theta - \beta iu + \frac{u^2}{2}} \right)$, $\text{Log}(z) = \ln|z| + i\text{Arg}(z)$, the $\text{Arg}(z)$ being chosen in $]-\Pi, \Pi]$, being the complexe logarithm. After derivation it follows $\Psi'(0) = i(\mu - \delta\beta)$ and $\Psi''(0) = \frac{\delta}{\theta}(\theta^2 - \beta^2)$, which yields $\alpha \equiv \frac{\mu - \delta\beta}{\theta^2 - \beta^2} \frac{\theta}{\delta}$.

3.3.3 Wiener integrals of Lévy processes

We take $X_t = \int_0^t \gamma_s d\Lambda_s$, where Λ is a square integrable Lévy process as in Section 3.3.2 with $\Lambda_0 = 0$. Then, $\int_0^T \gamma_s d\Lambda_s$ is well-defined at least when $\gamma \in \mathcal{L}^\infty([0, T])$. It is then possible to calculate the characteristic function and the cumulative function of $\int_0^T \gamma_s d\Lambda_s$. Let $(t, z) \mapsto t\Psi_\Lambda(z)$, denoting the log-characteristic function of Λ .

Lemma 3.12. *Let $\gamma : [0, T] \rightarrow \mathbb{R}$ be a Borel bounded function. The log-characteristic function of X_t is such that for all $u \in \mathbb{R}$, $\Psi_{X_t}(u) = \int_0^t \Psi_\Lambda(u\gamma_s) ds$, where $\mathbb{E}[\exp(iuX_t)] = \exp(\Psi_{X_t}(u))$. In particular, for every $t \in [0, T]$, $u \mapsto \Psi_{X_t}(u)$ is of class C^2 and so X_t is square integrable for any $t \in [0, T]$.*

Proof. Suppose first that γ is continuous, then $\int_0^T \gamma_s d\Lambda_s$ is the limit in probability of $\sum_{j=0}^{p-1} \gamma_{t_j} (\Lambda_{t_{j+1}} - \Lambda_{t_j})$ where $0 = t_0 < t_1 < \dots < t_p = T$ is a subdivision of $[0, T]$ whose mesh converges to zero. Using the independence of the increments, we have

$$\begin{aligned} \mathbb{E} \left[\exp \left\{ i \sum_{j=0}^{p-1} \gamma_{t_j} (\Lambda_{t_{j+1}} - \Lambda_{t_j}) \right\} \right] &= \prod_{j=0}^{p-1} \mathbb{E} [\exp \{ i \gamma_{t_j} (\Lambda_{t_{j+1}} - \Lambda_{t_j}) \}] = \prod_{j=0}^{p-1} \exp \{ \Psi_\Lambda(\gamma_{t_j})(t_{j+1} - t_j) \}, \\ &= \exp \left\{ \sum_{j=0}^{p-1} (t_{j+1} - t_j) \Psi_\Lambda(\gamma_{t_j}) \right\}. \end{aligned}$$

This converges to $\exp \left(\int_0^T \Psi_\Lambda(\gamma_s) ds \right)$, when the mesh of the subdivision goes to zero.

Suppose now that γ is only bounded and consider, using convolution, a sequence γ_n of continuous functions, such that $\gamma_n \rightarrow \gamma$ a.e. and $\sup_{t \in [0, T]} |\gamma_n(t)| \leq \sup_{t \in [0, T]} |\gamma(t)|$. We have proved that

$$\mathbb{E} \left[\exp \left(i \int_0^T \gamma_n(s) d\Lambda_s \right) \right] = \exp \left(\int_0^T \Psi_\Lambda(\gamma_n(s)) ds \right). \quad (3.13)$$

Now, Ψ_Λ is continuous therefore bounded, so Lebesgue dominated convergence and continuity of stochastic integral imply the statement. \square

Remark 3.13. 1. A similar statement was written with respect to the log cumulant generating function, see [4].
2. The proof works also when Λ has no moment condition and γ is a continuous function with bounded variation. Stochastic integrals are then defined using integration by parts.

Since Ψ_Λ is of class C^2 we have, $\Psi'_t(u) = \int_0^t \Psi'_\Lambda(u\gamma_s) \gamma_s ds$, and $\Psi''_t(u) = \int_0^t \Psi''_\Lambda(u\gamma_s) \gamma_s^2 ds$. So

$$\Psi'_t(0) = \Psi'_\Lambda(0) \int_0^t \gamma_s ds, \quad \Psi''_t(0) = \Psi''_\Lambda(0) \int_0^t \gamma_s^2 ds \quad \text{and} \quad \alpha_t = i \frac{\Psi'_\Lambda(0)}{\Psi''_\Lambda(0)} \frac{1_{\{\gamma_t \neq 0\}}}{\gamma_t}. \quad (3.14)$$

Remark 3.14. 1. $\text{Var}(X_T) = -\Psi''_\Lambda(0) \int_0^T \gamma_s^2 ds$.

2. If $\Psi''_\Lambda(0) = 0$ then $\text{Var}(X_T) = 0$ and so $\text{Var}(X_t) = 0, \forall t \in [0, T]$ and so X is deterministic. Consequently Condition (SC) is only verified if X vanishes identically because of Proposition 3.9.

Proposition 3.15. Condition (SC) is always verified if $\Psi''_\Lambda(0) \neq 0$.

Proof. We take into account item 2. of Proposition 3.7. Let $0 < s < t \leq T$, (3.14) implies

$$\Psi'_t(0) - \Psi'_s(0) = \Psi'_\Lambda(0) \int_s^t \gamma_r dr = \int_s^t \frac{\Psi'_\Lambda(0)}{\gamma_r} 1_{\{\gamma_r \neq 0\}} \gamma_r^2 dr = \int_s^t (-i\alpha_r) d(-\Psi''_r(0))$$

where $\alpha_r = \frac{\Psi'_\Lambda(0)}{\Psi''_\Lambda(0)} \frac{i}{\gamma_r} 1_{\{\gamma_r \neq 0\}}$. This shows the first point of (3.9). In particular $\left| \frac{d\Psi'_t(0)}{d\Psi''_t(0)} \right| = \frac{|\Psi'_\Lambda(0)|}{-\Psi''_\Lambda(0)} \frac{1}{\gamma_t}$. The second point of (3.9) follows because $\int_0^T |i\alpha_r|^2 d\left(\Psi''_r(0)\right) = T \frac{|\Psi'_\Lambda(0)|^2}{(-\Psi''_\Lambda(0))^2} < \infty$. \square

4 Explicit Föllmer-Schweizer decomposition in the PII case

Let $X = (X_t)_{t \in [0, T]}$ be a semimartingale (measurable process) with independent increments with log-characteristic function $(t, u) \mapsto \Psi_t(u)$. We assume that $(X_t)_{t \in [0, T]}$ is square integrable. In this section, we first evaluate an explicit Kunita-Watanabe decomposition of a random variable H w.r.t. the martingale part M of X . Later, we obtain the decomposition with respect to X . Before doing so, it is useful to introduce in the following preliminary subsection an *expectation* operator and a *derivative* operator related to X .

From now on we will suppose the validity of the (SC) condition.

4.1 On some expectation and derivative operators

We first introduce the *expectation* operator related to X . For $0 \leq t \leq T$, let $\epsilon_{t, T}^X$ denote the complex valued function defined for all $u \in \mathbb{R}$ by

$$\epsilon_{t, T}^X(u) := \exp(\Psi_T(u) - \Psi_t(u)). \quad (4.1)$$

In the sequel, to simplify notations, we will write $\epsilon_{t, T}$ instead of $\epsilon_{t, T}^X$.

We observe that the function $(u, t) \mapsto \epsilon_{t, T}(u)$ and $(u, t) \mapsto \epsilon_{t, T}^2(u)$ are uniformly bounded because the characteristic function is bounded. The lemma below shows that the function $\epsilon_{t, T}$ is closely related to the conditional expectation.

Lemma 4.1. *Let $H = f(X_T)$ where f is given as a Fourier transform, $f(x) := \hat{\mu}(x) := \int_{\mathbb{R}} e^{iux} \mu(du)$, of a (finite) complex measure μ defined on \mathbb{R} .*

Then, for all $t \in [0, T]$, $\mathbb{E}[f(X_T) | \mathcal{F}_t] = e_{t, T}(X_t)$ where for all $x \in \mathbb{R}$,

$$e_{t, T}(x) := \widehat{\epsilon_{t, T} \mu}(x) = \int_{\mathbb{R}} e^{iux} \epsilon_{t, T}(u) \mu(du).$$

Proof. First, we easily check that $\int_{\mathbb{R}} \epsilon_{t, T}(u) \mu(du) < \infty$, since μ is supposed to be a finite measure.

Now, let us consider the conditional expectation $\mathbb{E}[f(X_T) | \mathcal{F}_t]$. By Fubini's theorem,

$$\mathbb{E}[f(X_T) | \mathcal{F}_t] = \mathbb{E} \left[\int_{\mathbb{R}} e^{iuX_T} \mu(du) | \mathcal{F}_t \right] = \int_{\mathbb{R}} \mu(du) \mathbb{E}[e^{iuX_T} | \mathcal{F}_t].$$

Finally, remark that by the independent increments property of X , we obtain

$$\mathbb{E}[e^{iuX_T} | \mathcal{F}_t] = \mathbb{E} \left[e^{iu(X_T - X_t)} e^{iuX_t} | \mathcal{F}_t \right] = \exp(\Psi_T(u) - \Psi_t(u)) e^{iuX_t}, \quad \text{for all } u \in \mathbb{R}.$$

\square

Now let us introduce the *derivative* operator related to the PII X . Let δ_t^X denote the complex valued function defined for all $u \in \mathbb{R}$ by the Radon-Nykodim derivative

$$\delta_t^X(u) := i \frac{d(\Psi'_t(u) - \Psi'_t(0))}{d\Psi''_t(0)}, \quad (4.2)$$

which is well-defined by Corollary 3.11. In the sequel, to simplify notations, we will write δ_t instead of δ_t^X . By (3.7) in Remark 3.3 4., we obtain

$$\delta_t(u) = iuc_t + \int_{\mathbb{R}} x(e^{iux} - 1)F_t(dx) . \quad (4.3)$$

The lemma below shows that the function δ_t is closely related to the Malliavin derivative in the sense of [3].

Lemma 4.2. *Let η be a finite complex measure defined on \mathbb{R} with a finite first order moment and g its Fourier transform, i.e. the complex-valued function such that for all $x \in \mathbb{R}$, $g(x) = \hat{\eta}(x) := \int_{\mathbb{R}} e^{iux} \eta(du)$.*

1. *g is differentiable with bounded derivative;*
2. *$\delta_t(u)\eta(du)$ is a finite complex measure.*
3. *For all $x \in \mathbb{R}$,*

$$\widehat{\delta_t \eta}(x) := \int_{\mathbb{R}} e^{iux} \delta_t(u) \eta(du) = c_t g'(x) + \int_{\mathbb{R}} (g(x+y) - g(x)) y F_t(dy) . \quad (4.4)$$

Proof. Item 1. is obvious. We prove item 2. i.e. that $\int_{\mathbb{R}} |\delta_t \eta|(du) < \infty$. For this, notice that the following upper bound holds for all $u, x \in \mathbb{R}$,

$$|x(e^{iux} - 1)| = 2|x| \left| \sin \frac{ux}{2} \right| \leq 2(|u| \vee 1)(x^2 \wedge |x|) . \quad (4.5)$$

Now, using the expression (4.3) of δ_t yields

$$|\delta_t(u)| \leq \sqrt{2} \left[c_t |u| + 2(1 + |u|) \int_{\mathbb{R}} x^2 F_t(dx) \right] \leq 2\sqrt{2}(1 + |u|) , \quad (4.6)$$

because by point 2. of Corollary 3.11, $c_t \leq 1$ and $\int_{\mathbb{R}} x^2 F_t(dx) \leq 1$ $d(-\Psi_t''(0))$ -a.e. Finally, (4.6) and the fact that η is supposed to have a finite first order moment imply the result.

We go on with the proof of point 3. Now we can consider the Fourier transform $\widehat{\delta_t \eta}$. Using Fubini's theorem and (4.6), we obtain the following expression

$$\begin{aligned} \widehat{\delta_t \eta}(x) &= \int_{\mathbb{R}} \delta_t(u) \eta(du) e^{iux} \\ &= c_t \int_{\mathbb{R}} iu \eta(du) e^{iux} + \int_{\mathbb{R}} \left(\int_{\mathbb{R}} \eta(du) e^{iu(x+y)} - \int_{\mathbb{R}} \eta(du) e^{iux} \right) y F_t(dy) . \end{aligned}$$

□

We are now in the position to state an explicit expression for the Kunita-Watanabe decomposition of some random variables of the form $H = f(X_T)$. To be more specific, we consider a random variable which is given as a Fourier transform of X_T ,

$$H = f(X_T) \quad \text{with} \quad f(x) = \hat{\mu}(x) = \int_{\mathbb{R}} e^{iux} \mu(du) , \quad \text{for all } x \in \mathbb{R} \quad (4.7)$$

for some finite complex signed measure μ .

4.2 Explicit elementary Kunita-Watanabe decomposition

By Proposition 3.7, X admits the following semimartingale decomposition, $X_t = A_t + M_t$, where

$$A_t = -i\Psi'_t(0) \quad \text{and} \quad \langle M \rangle_t = -\Psi''_t(0). \quad (4.8)$$

Proposition 4.3. *Let $H = f(X_T)$ where f is of the form (4.7). Then, H admits the decomposition*

$$H = \mathbb{E}[H] + \int_0^T Z_t dM_t + O_T,$$

where O is a square integrable (\mathcal{F}_t) -martingale such that $\langle O, M \rangle = 0$ and

$$V_t := \mathbb{E}[H | \mathcal{F}_t] = e_{t,T}(X_t), \quad \text{and} \quad Z_t = d_{t,T}(X_{t-}),$$

where the complex valued functions $e_{t,T}$ and $d_{t,T}$ are defined for all $x \in \mathbb{R}$ by

$$e_{t,T}(x) := \widehat{\epsilon_{t,T}\mu}(x) = \int_{\mathbb{R}} \epsilon_{t,T}(u) \mu(du) e^{iux} \quad \text{and} \quad d_{t,T}(x) := \widehat{\delta_t \epsilon_{t,T}\mu}(x) = \int_{\mathbb{R}} \delta_t(u) \epsilon_{t,T}(u) \mu(du) e^{iux}, \quad (4.9)$$

with $\epsilon_{t,T}$ being defined in (4.1) and δ_t being defined in (4.2). Moreover, $\mathbb{E}[\int_0^T Z_s^2 d\langle M \rangle_s] < \infty$.

In particular, $V_0 = \mathbb{E}[H]$.

Remark 4.4. We remark that the (SC) condition is not a restriction when X is a martingale, since it is obviously fulfilled. This would correspond to the classical Kunita-Watanabe statement.

Remark 4.5. In [3], they obtain a similar decomposition valid for a different class of random variables. On one hand their class is more general, allowing for path dependent payoffs, on the other hand it requires some stronger regularity assumptions since H is supposed to be in the Malliavin-Sobolev space $\mathbb{D}^{1,2}$. In our case, their regularity assumption on the payoff function could be relaxed by applying the derivative operator δ_t after applying the expectation operator $\epsilon_{t,T}$ whereas in [3], they take the conditional expectation of the payoff Malliavin derivative.

This trick of switching the conditional expectation and the differentiation is also implicitly used in the approach developed in [16] or similarly in [7]. Their approach relies on the application of Ito's lemma on the conditional expectation $\mathbb{E}[H | \mathcal{F}_t]$ and therefore requires some regularity conditions. Basically the conditional expectation should be once differentiable w.r.t. the time variable and twice differentiable w.r.t. the space variable with continuous partial derivatives. On the other hand, their method is valid for a large class of martingale processes X .

Besides, our approach only relies on the martingale property of $(e^{iuX_t - \Psi_t(u)})_{0 \leq t \leq T}$. Hence, X is not required to be martingale as in [3], [16] or [7] and no specific regularity assumption on the payoff function or on the conditional expectation are required. Our approach is unfortunately restricted to additive processes. However, this specific setting allows to go one step further in providing an explicit expression for both the Follmer-Schweizer decomposition and the variance optimal strategy, as we will see below. Moreover, the expression of the Kunita-Watanabe decomposition derived in this specific case is quasi-explicit involving a simple Fourier transform.

If $\epsilon_{t,T}\mu$ admits a first order moment, then taking $\eta = \epsilon_{t,T}\mu$, in Lemma 4.2, the conditional expectation function $e_{t,T}$ is differentiable w.r.t. the variable x and we obtain

$$d_{t,T}(x) := \widehat{\delta_t \epsilon_{t,T}\mu}(x) = \int_{\mathbb{R}} \delta_t(u) \epsilon_{t,T}(u) \mu(du) e^{iux} = c_t e'_{t,T}(x) + \int_{\mathbb{R}} (e_{t,T}(x+y) - e_{t,T}(x)) y F_t(dy). \quad (4.10)$$

The following lemma gives a condition on characteristics c_t and F_t ensuring the differentiability of $e_{t,T}$.

Lemma 4.6. Let X be a PII process with finite second order moments such that there exist positive reals $\beta \in (0, 2)$ and α verifying

$$\inf_{t \in [0, T)} \left(c_t + \int_{|x| \leq |u|^{-1}} x^2 F_t(dx) \right) \geq \alpha |u|^{-2+\beta}, \quad \text{when } |u| \rightarrow \infty. \quad (4.11)$$

Let μ be a finite complex measure defined on \mathbb{R} and f its Fourier transform such that for all $x \in \mathbb{R}$, $f(x) = \hat{\mu}(x)$. Then, for all $t \in [0, T)$, $\epsilon_{t,T}\mu$ is a finite complex measure with finite moments of all orders and all the derivatives of all orders of $x \mapsto \epsilon_{t,T}(x) := \widehat{\epsilon_{t,T}\mu}(x)$ are well-defined and bounded.

Remark 4.7. When X is a Lévy process, Assumption (4.11) implies the Kallenberg condition stated in [18] ensuring the existence of a transition density for a Lévy process X .

Proof of Lemma 4.6. We prove that $\int_{\mathbb{R}} u^p \epsilon_{t,T}(u) \mu(du) < \infty$, for any nonnegative integer p . For this, we recall that Remark 3.3 together with the lines below Proposition 3.10 say that for all $u \in \mathbb{R}$ and $t \in [0, T]$

$$|\epsilon_{t,T}(u)| = \exp \left\{ -\frac{u^2}{2} \int_t^T c_s d(-\Psi_s''(0)) \right\} \left| \exp \left\{ \int_t^T \int_{\mathbb{R}} (e^{iux} - 1 - iux \mathbf{1}_{|x| \leq 1}) F_s(dx) d(-\Psi_s''(0)) \right\} \right|. \quad (4.12)$$

Consider now the second exponential term on the right-hand side of the above equality; it gives

$$\begin{aligned} \left| \exp \left\{ \int_t^T \int_{\mathbb{R}} (e^{iux} - 1 - iux \mathbf{1}_{|x| \leq 1}) F_s(dx) d(-\Psi_s''(0)) \right\} \right| &\leq \exp \left\{ -2 \int_t^T \int_{\mathbb{R}} \left(\sin \frac{ux}{2} \right)^2 F_s(dx) d(-\Psi_s''(0)) \right\} \\ &\leq \exp \left\{ -2 \int_t^T \int_{|x| \leq \frac{\pi}{|u|}} \left(\sin \frac{ux}{2} \right)^2 F_s(dx) d(-\Psi_s''(0)) \right\} \\ &\leq \exp \left\{ -2 \int_t^T \left(\frac{u}{\pi} \right)^2 \int_{|x| \leq \frac{\pi}{|u|}} x^2 F_s(dx) d(-\Psi_s''(0)) \right\}. \end{aligned} \quad (4.13)$$

Hence, we conclude that for all $u \in \mathbb{R}$,

$$|\epsilon_{t,T}(u)| \leq \exp \left\{ -\frac{u^2}{2} \int_t^T \left[c_s + \frac{4}{\pi^2} \int_{|x| \leq \frac{\pi}{|u|}} x^2 F_s(dx) \right] d(-\Psi_s''(0)) \right\}. \quad (4.14)$$

Then by Assumption (4.11), there exists two positive reals α and β such that

$$|\epsilon_{t,T}(u)| \leq \exp \left\{ -\alpha \left(\frac{u}{\pi} \right)^\beta (\Psi_t''(0) - \Psi_T''(0)) \right\}, \quad \text{as } |u| \rightarrow \infty.$$

Finally, we can conclude that under Assumption (4.11), for any nonnegative integer p , the complex measure $u^p \epsilon_{t,T}\mu$ is finite with a bounded Fourier transform $g_{t,T}^{(p)}$, the p order derivative of $\epsilon_{t,T}$. \square

Remark 4.8. Notice that the explicit expression of the Kunita-Watanabe decomposition obtained in the case of an additive process can be used to derive an explicit expression in the case where X is an Ornstein-Uhlenbeck process. Indeed if we consider

$$X_t = e^{-\alpha t} \tilde{X}_t, \quad (4.15)$$

for a given positive real $\alpha \in \mathbb{R}$ and additive process \tilde{X} . Consider a function f satisfying condition (4.7). We define now $\tilde{f} : [0, T] \times \mathbb{R} \rightarrow \mathbb{R}$ by $\tilde{f}(t, \tilde{x}) = f(e^{-\alpha t} \tilde{x})$, for all $\tilde{x} \in \mathbb{R}$, so that $f(X_T) = \tilde{f}(T, \tilde{X}_T)$. Then by application of Proposition 4.3, we get $f(X_T) = \mathbb{E}[f(X_T)] + \int_0^T \tilde{Z}_s^T d\tilde{M}_s + O_T$, where \tilde{M} is the martingale part of \tilde{X} . Now $dX_t = -\alpha e^{-\alpha t} \tilde{X}_t dt + e^{-\alpha t} d\tilde{M}_t - i e^{-\alpha t} d\Psi_t'(0)$, where Ψ is the log-characteristic function of \tilde{X} . By uniqueness of

the Doob-Meyer decomposition of the special semimartingale X , the martingale part of X is $M_t = \int_0^t e^{-\alpha s} d\tilde{M}_s$ and finally we deduce the Kunita-Watanabe decomposition

$$f(X_T) = \mathbb{E}[f(X_T)] + \int_0^T Z_s dM_s + O_T, \quad \text{with } Z_t = e^{\alpha t} \tilde{Z}_t^T. \quad (4.16)$$

This can be easily generalized when αt is replaced by $\alpha(t)$ a bounded deterministic function of t .

Proof of Proposition 4.3. Lemma 4.1 says that $V_t := \mathbb{E}[H|\mathcal{F}_t] = \int V_t(u) d\mu(u)$ where

$$V_t(u) = \epsilon_{t,T}(u) e^{iuX_t} = \mathbb{E}[\exp(iuX_T) | \mathcal{F}_t]. \quad (4.17)$$

Having observed that $|\epsilon_{t,T}(u)| \leq 1$, for all $u \in \mathbb{R}$, we get

$$\sup_{t \leq T, u \in \mathbb{R}} \mathbb{E}[|V_t(u)|^2] \leq 1. \quad (4.18)$$

This implies that V is an (\mathcal{F}_t) -square integrable martingale since μ is finite. We define

$$Z_t = \int_{\mathbb{R}} Z_t(u) d\mu(u), \quad (4.19)$$

where

$$Z_t(u) = \delta_{t,T}(u) \epsilon_{t,T}(u) e^{iuX_{t-}}. \quad (4.20)$$

In the second part of this proof, we will show that

$$\mathbb{E} \left[\int_0^T |Z_t(u)|^2 d(-\Psi_t''(0)) \right] \leq 2. \quad (4.21)$$

This implies in particular that process Z in (4.19) is well defined and $\mathbb{E} \left[\int_0^T |Z_s|^2 d\langle M \rangle_s \right] < \infty$. We define

$$O_t := V_t - V_0 - \int_0^t Z_s dM_s. \quad (4.22)$$

By additivity O is an (\mathcal{F}_t) -square integrable martingale. It remains to prove that $\langle O, M \rangle = 0$. For this, we will show that

$$\langle V, M \rangle = \int_0^t Z_s d\langle M \rangle_s,$$

which will follow from the fact that $V_t M_t - \int_0^t Z_s d\langle M \rangle_s$ is an (\mathcal{F}_t) -martingale. In order to establish the latter, we prove that for every $0 < r < t$,

$$\mathbb{E} \left[\left(V_t M_t - V_r M_r - \int_r^t Z_s d\langle M \rangle_s \right) R_r \right] = 0 \quad (4.23)$$

for every bounded (\mathcal{F}_r) -measurable variable R_r . Taking into account (4.18) and (4.21), by Fubini's theorem, the left hand side of (4.23) equals

$$\int d\mu(u) \mathbb{E} \left[\left(V_t(u) M_t - V_r(u) M_r - \int_r^t Z_s(u) d(-\Psi_s''(0)) \right) R_r \right]. \quad (4.24)$$

It remains now to show that the expectation in (4.24) vanishes for $d\mu(u)$ almost all u . Below, we will show that

$$V_t(u) M_t - \int_0^t Z_s(u) d(-\Psi_s''(0))$$

is an (\mathcal{F}_t) -martingale, for $d\mu(u)$ almost all u . This implies that (4.24) is zero.

We evaluate $\mathbb{E}[V_t M_t | \mathcal{F}_s]$. Since V and M are (\mathcal{F}_t) -martingales, using the property of independent increments we get

$$\begin{aligned} \mathbb{E}[V_t(u) M_t | \mathcal{F}_s] &= \mathbb{E}[V_t(u) M_s | \mathcal{F}_s] + \mathbb{E}[V_t(u) (M_t - M_s) | \mathcal{F}_s], \\ &= M_s V_s(u) + V_s(u) \mathbb{E}[\exp\{iu(X_t - X_s) - (\Psi_t(u) - \Psi_s(u))\} (M_t - M_s)], \\ &= M_s V_s(u) + V_s(u) e^{-(\Psi_t(u) - \Psi_s(u))} \mathbb{E}[e^{iu(X_t - X_s)} (M_t - M_s)]. \end{aligned}$$

Consider now the expectation on the right hand side of the above equality:

$$\begin{aligned} \mathbb{E}[e^{iu(X_t - X_s)} (M_t - M_s)] &= \mathbb{E}[e^{iu(X_t - X_s)} (X_t - X_s)] + \mathbb{E}[e^{iu(X_t - X_s)} i(\Psi'_t(0) - \Psi'_s(0))], \\ &= -i \frac{\partial}{\partial u} \mathbb{E}[e^{iu(X_t - X_s)}] + i(\Psi'_t(0) - \Psi'_s(0)) \mathbb{E}[e^{iu(X_t - X_s)}], \\ &= -i e^{\Psi_t(u) - \Psi_s(u)} (\Psi'_t(u) - \Psi'_s(u)) + i(\Psi'_t(0) - \Psi'_s(0)) e^{\Psi_t(u) - \Psi_s(u)}. \end{aligned}$$

Consequently,

$$\begin{aligned} \mathbb{E}[V_t(u) M_t | \mathcal{F}_s] &= M_s V_s(u) - i V_s(u) (\Psi'_t(u) - \Psi'_s(u)) + i V_s(u) (\Psi'_t(0) - \Psi'_s(0)) \\ &= M_s V_s(u) - i V_s(u) (\Psi'_t(u) - \Psi'_t(0) - (\Psi'_s(u) - \Psi'_s(0))). \end{aligned}$$

This implies that $(V_t(u) M_t + i V_t(u) (\Psi'_t(u) - \Psi'_t(0)))_t$ is an (\mathcal{F}_t) -martingale. Then by integration by parts,

$$V_t(u) (\Psi'_t(u) - \Psi'_t(0)) = \int_0^t V_s(u) d(\Psi'_s(u) - \Psi'_s(0)) + \int_0^t (\Psi'_s(u) - \Psi'_s(0)) dV_s(u).$$

The proof is concluded once we have shown (4.21).

By Cauchy-Schwarz inequality

$$\begin{aligned} \mathbb{E} \left(\int_0^T Z_s^2 d(-\Psi''_s(0)) \right) &= \mathbb{E} \left(\int_0^T \left| \int_{\mathbb{R}} \delta_s(u) \epsilon_{s,T}(u) e^{iuX_s} \mu(du) \right|^2 d(-\Psi''_s(0)) \right) \\ &\leq |\mu|(\mathbb{R}) \int_0^T \int_{\mathbb{R}} |\delta_s(u) \epsilon_{s,T}(u)|^2 |\mu|(du) d(-\Psi''_s(0)) \\ &= |\mu|(\mathbb{R}) \int_{\mathbb{R}} |\mu|(du) \int_0^T |\delta_s(u) \epsilon_{s,T}(u)|^2 d(-\Psi''_s(0)). \end{aligned}$$

Let us consider now, for a given real u the integral w.r.t. to the time parameter in the right-hand side of the above inequality. Using inequalities (4.13) and (4.12), we obtain

$$|\epsilon_{s,T}(u)|^2 \leq -\frac{1}{2} \exp \left\{ \int_s^T - \left[u^2 c_r + 4 \int_{\mathbb{R}} (\sin \frac{ux}{2})^2 F_r(dx) \right] d(-\Psi''_r(0)) \right\}.$$

On the other hand, by (4.3) and (4.5) we have

$$\begin{aligned} |\delta_s(u)|^2 &\leq 2 \left[u^2 c_s^2 + 4 \left(\int_{\mathbb{R}} |x \sin \frac{ux}{2}| F_s(dx) \right)^2 \right] \\ &\leq 2 \left[u^2 c_s^2 + 4 \int_{\mathbb{R}} x^2 F_s(dx) \int_{\mathbb{R}} (\sin \frac{ux}{2})^2 F_s(dx) \right]. \end{aligned}$$

Since $c_s \leq 1$, $\int_{\mathbb{R}} x^2 F_s(dx) \leq 1$, by item 2. of Corollary 3.11, we finally get

$$|\delta_s(u)|^2 \leq 2\gamma_s(u) \quad (4.25)$$

where

$$\gamma_s(u) = u^2 c_s + 4 \int_{\mathbb{R}} \left(\sin \frac{ux}{2}\right)^2 F_s(dx). \quad (4.26)$$

Then

$$\begin{aligned} \int_0^T |\delta_s(u) \epsilon_{s,T}(u)|^2 d(-\Psi_s''(0)) &\leq 2 \int_0^T \gamma_s(u) \exp \left\{ \int_s^T -\gamma_r(u) d(-\Psi_r''(0)) \right\} d(-\Psi_s''(0)) \\ &= 2 \left(1 - \exp \left\{ \int_0^T -\gamma_r(u) d(-\Psi_r''(0)) \right\} \right) \leq 2. \end{aligned} \quad (4.27)$$

□

Example 4.9. We take $X = M = W$ the classical Wiener process with canonical filtration (\mathcal{F}_t) . We have $\Psi_s(u) = -\frac{u^2 s}{2}$ so that $\Psi_s'(u) = -us$ and $\Psi_s''(u) = -s$. So $Z_s(u) = iuV_s(u)$. We recall that $V_s = \mathbb{E}[\exp(iuW_T)|\mathcal{F}_s] = \exp(iuW_s) \exp(-u^2 \frac{T-s}{2})$. In particular, $V_0 = \exp(-\frac{u^2 T}{2})$ and so $\exp(iuW_T) = i \int_0^T u \exp(iuW_s) \exp(-u^2 \frac{T-s}{2}) dW_s + \exp(-\frac{u^2 T}{2})$. In fact that expression is classical and it can be derived from Clark-Ocone formula. In fact, if D is the usual Malliavin derivative then $\mathbb{E}(D_t \exp(iuW_T)|\mathcal{F}_t) = iu \exp(iuW_s - \frac{u^2}{2}(T-s))$.

4.3 Explicit Föllmer-Schweizer decomposition

We are now able to evaluate the FS decomposition of $H = f(X_T)$ where f is given by (4.7). First, we state the following lemma.

Lemma 4.10. For all $s, t \in [0, T)$,

$$\int_s^t \operatorname{Re}(i\delta_r(u) d\Psi_r'(0)) \leq K_T + \int_s^t \int_{\mathbb{R}} \left(\sin \frac{ux}{2}\right)^2 F_r(dx) d(-\Psi_r''(0)), \quad (4.28)$$

where the process K was defined in Definition 2.3.

Proof. Using (4.3) and (3.7), with a slight abuse of notation, it follows

$$\begin{aligned} \operatorname{Re}(i\delta_r(u) d\Psi_r'(0)) &= - \left(b_r + \int_{|x|>1} x F_r(dx) \right) \int_{\mathbb{R}} (x(\cos(ux) - 1) F_r(dx)) d(-\Psi_r''(0)) \\ &= 2 \left(b_r + \int_{|x|>1} x F_r(dx) \right) \int_{\mathbb{R}} \left(x \left(\sin \frac{ux}{2} \right)^2 F_r(dx) \right) d(-\Psi_r''(0)) \\ &\leq \left[\left(b_r + \int_{|x|>1} x F_r(dx) \right)^2 + \left(\int_{\mathbb{R}} x \left(\sin \frac{ux}{2} \right)^2 F_r(dx) \right)^2 \right] d(-\Psi_r''(0)). \end{aligned}$$

Indeed, by the (SC) condition using Proposition 3.7 and Corollary 3.11 1., we obtain that for all $t \in [0, T)$,

$$\int_s^t \left(b_r + \int_{|x|>1} x F_r(dx) \right)^2 d(-\Psi_r''(0)) = \int_s^t \left| \frac{d\Psi_r'(0)}{d(-\Psi_r''(0))} \right|^2 d(-\Psi_r''(0)) = K_t - K_s \leq K_T,$$

which is a deterministic bound. Finally, recalling that $\int_{\mathbb{R}} x^2 F_r(dx) \leq 1$ by Corollary 3.11 2., Cauchy-Schwarz inequality implies

$$\left(\int_{\mathbb{R}} x \left(\sin \frac{ux}{2} \right)^2 F_r(dx) \right)^2 \leq \int_{\mathbb{R}} \left(\sin \frac{ux}{2} \right)^2 F_r(dx).$$

□

Theorem 4.11. *The FS decomposition of $H = f(X_T)$ where f satisfies (4.7) is the following*

$$H_t = H_0 + \int_0^t \xi_s dX_s + L_t \quad \text{with} \quad H_T = H, \quad (4.29)$$

where

$$\xi_t = k_{t,T}(X_{t-}), \quad \text{with} \quad k_{t,T}(x) = \int_{\mathbb{R}} e^{i \int_t^T \delta_s(u) d\Psi'_s(0)} \delta_t(u) \epsilon_{t,T}(u) e^{iux} \mu(du), \quad (4.30)$$

and

$$H_t = h_{t,T}(X_t), \quad \text{with} \quad h_{t,T}(x) = \int_{\mathbb{R}} e^{i \int_t^T \delta_s(u) d\Psi'_s(0)} \epsilon_{t,T}(u) e^{iux} \mu(du), \quad (4.31)$$

with $\epsilon_{t,T}$ defined in (4.1) and δ_t defined in (4.2).

Proof. Let us introduce the following notations, which will correspond to the expression (4.31) for H_t and (4.30) for ξ_t in the case where $\mu = \delta_u$ for a given real u :

$$H_t(u) := e^{i \int_t^T \delta_s(u) d\Psi'_s(0)} \epsilon_{t,T}(u) e^{iux_t} \quad \text{and} \quad \xi_t(u) := e^{i \int_t^T \delta_s(u) d\Psi'_s(0)} \delta_t(u) \epsilon_{t,T}(u) e^{iux_{t-}}. \quad (4.32)$$

1. We first introduce the process H . Taking into account Lemma 4.10 together with inequalities (4.13) and (4.12), $|H(u)|$ is uniformly bounded in u and t . Indeed

$$\begin{aligned} |H_t(u)| &= |e^{i \int_t^T \delta_r(u) d\Psi'_r(0)}| |\epsilon_{t,T}(u)| \\ &\leq \exp \left\{ K_T + \int_t^T \int_{\mathbb{R}} \left(\sin \frac{ux}{2} \right)^2 F_r(dx) d(-\Psi''_r(0)) \right\} \\ &\quad \exp \left\{ \int_t^T -\frac{1}{2} \left[u^2 c_r + 4 \int_{\mathbb{R}} \left(\sin \frac{ux}{2} \right)^2 F_r(dx) \right] d(-\Psi''_r(0)) \right\} \\ &\leq \exp(K_T) \exp \left\{ -\frac{1}{4} \int_t^T \gamma_r(u) d(-\Psi''_r(0)) \right\} \\ &\leq \exp(K_T), \end{aligned} \quad (4.33)$$

where γ was defined in (4.26).

By Fubini's,

$$H_t = \int_{\mathbb{R}} H_t(u) d\mu(u) \quad (4.34)$$

is well-defined and it equals the expression in (4.31). We prove now that ξ defined in (4.30) is a well-defined square integrable process. Using the above bounds (4.33) and (4.25) we obtain

$$\begin{aligned} |\xi_t(u)|^2 &\leq |\delta_t(u)|^2 |H_t(u)|^2 \\ &\leq 2\gamma_t(u) \exp(2K_T) \exp \left\{ -\frac{1}{2} \int_t^T \gamma_r(u) d(-\Psi''_r(0)) \right\}, \end{aligned} \quad (4.35)$$

which finally implies

$$\mathbb{E} \left(\int_0^T |\xi_t(u)|^2 d\langle M \rangle_t \right) \leq 4 \left(1 - \exp \left\{ \int_t^T \gamma_r(u) d(-\Psi''_r(0)) \right\} \right) \leq 4. \quad (4.36)$$

Hence, $\xi(u) \in \Theta := L^2(M)$ for any $u \in \mathbb{R}$. (4.36), (4.27) yield

$$\int_{\mathbb{R}} d\mu(u) \int_0^T \mathbb{E}(|\xi_t(u)|^2) d\langle M \rangle_t < \infty. \quad (4.37)$$

The above upper bound implies that $\xi_t = \int_{\mathbb{R}} \xi_t(u) d\mu(u)$ is well-defined and $\xi \in \Theta$ and it equals the expression (4.30). Consequently $\xi \in L^2(M) = \Theta$ and using stochastic and classical Fubini's we get

$$\int_0^t \xi_s dX_s = \int_{\mathbb{R}} d\mu(u) \int_0^t \xi_s(u) dX_s. \quad (4.38)$$

2. We go on with the proof of Theorem 4.11 showing the following:
 - a) $L_t = H_t - H_0 - \int_0^t \xi_s dX_s$ is an eventually complex valued square integrable martingale;
 - b) $\langle L, M \rangle = 0$ where M is the martingale part of the special semimartingale X .
3. We first establish a) and b) for the case μ is the Dirac measure at some fixed $u \in \mathbb{R}$. We will show that

$$H_t(u) = H_0(u) + \int_0^t \xi_s(u) dX_s + L_t(u) \quad \text{with} \quad H_T(u) = \exp(iuX_T), \quad (4.39)$$

for fixed $u \in \mathbb{R}$ where $L(u)$ is a square integrable martingale and $\langle L(u), M \rangle = 0$. Notice that by relation (4.32), $H_t(u) = e^{i \int_0^t \delta_s(u) d\Psi'_s(0)} V_t(u)$ with $V_t(u) = e^{iuX_t} \epsilon_{t,T}(u)$, as introduced in (4.17). Integrating by parts, gives

$$H_t(u) = H_0(u) + \int_0^t e^{i \int_r^T \delta_s(u) d\Psi'_s(0)} V_r(u) (-i\delta_r(u) d\Psi'_r(0)) + \int_0^t e^{i \int_r^T \delta_s(u) d\Psi'_s(0)} dV_r(u). \quad (4.40)$$

We denote again by $Z(u)$ the expression provided by (4.20). We observe that

$$\xi_t(u) = e^{i \int_t^T \delta_s(u) d\Psi'_s(0)} Z_t(u).$$

We recall that

$$dV_r(u) = Z_r(u) dM_r + dO_r(u) = Z_r(u) (dX_r - dA_r) + dO_r(u), \quad (4.41)$$

where A is given by (4.8) and $O(u)$ is a square integrable martingale strongly orthogonal to M . Replacing (4.41) in (4.40) yields

$$\begin{aligned} H_t(u) &= H_0(u) + L_t(u) + \int_0^t e^{i \int_r^T \delta_s(u) d\Psi'_s(0)} Z_r(u) dX_r \\ &+ i \int_0^t e^{i \int_r^T \delta_s(u) d\Psi'_s(0)} Z_r(u) d\Psi'_r(0) - i \int_0^t e^{i \int_r^T \delta_s(u) d\Psi'_s(0)} V_r(u) \delta_r(u) d\Psi'_r(0) \\ &= H_0(u) + L_t(u) + \int_0^t \xi_s(u) dX_s, \end{aligned}$$

where

$$L_t(u) = \int_0^t e^{i \int_r^T \delta_s(u) d\Psi'_s(0)} dO_r(u). \quad (4.42)$$

$L(u)$ is a local martingale which is also a square integrable martingale because $\int_0^T e^{2\operatorname{Re}(i \int_t^T \delta_s(u) d\Psi'_s(0))} d\langle O \rangle_t$ is finite taking into account Lemma 4.10.

Since $O(u)$ is strongly orthogonal with respect to M , then $L(u)$ has the same property.

4. We treat now the general case discussing the points a) and b) in item 2. (4.38) and the Definition of H show that

$$L_t := H_t - H_0 - \int_0^t \xi_s dX_s \quad (4.43)$$

fulfills

$$\int_{\mathbb{R}} L_t(u) d\mu(u) = L_t, \quad (4.44)$$

for every $t \in [0, T]$. Let $0 \leq s < t \leq T$ and R_s a bounded \mathcal{F}_s -measurable random variable. Using (4.33), (4.36) and Cauchy-Schwarz we obtain

$$\begin{aligned} \mathbb{E}(|L_t(u)|^2) &= \mathbb{E}\left(\left|H_t(u) - H_0(u) - \int_0^t \xi_r(u) dX_r\right|^2\right) \\ &\leq 2\mathbb{E}(|H_t(u)|^2) + 4\mathbb{E}(|H_0(u)|^2) + 8\left(\mathbb{E}\left(\int_0^t \xi_r(u) dM_r\right)^2 + \mathbb{E}\left(\int_0^t \xi_r(u) \alpha_r d[M]_r\right)^2\right) \\ &\leq 2\mathbb{E}(|H_t(u)|^2) + 4\mathbb{E}(|H_0(u)|^2) + 8(1 + K_T) \mathbb{E}\left(\int_0^t |\xi_r(u)|^2 d(-\Psi_r''(0))\right) \\ &\leq 6\exp(2K_T) + 32(1 + K_T). \end{aligned} \quad (4.45)$$

- (a) By (4.45), we observe that $\mathbb{E}(\int_{\mathbb{R}} d\mu(u) |L_t(u)|) < \infty$. Fubini's, (4.44) and the fact that $L(u)$ is an (\mathcal{F}_t) -martingale give $\mathbb{E}[L_t R_s] = \mathbb{E}[L_s R_s]$. Therefore L is an (\mathcal{F}_t) -martingale. For every $t \in [0, T]$, (L_t) is a square integrable because of (4.43) and by additivity.
- (b) By item 3. $L(u)M$ is an (\mathcal{F}_t) -local martingale. Moreover $L(u)$ and M are square integrable martingales. By Cauchy-Schwarz and Doob inequalities, it follows that $E(\sup_{t \in [0, T]} |L_t(u) M_t|)$ is finite. Consequently $L(u)M$ is indeed an (\mathcal{F}_t) -martingale. It remains to show that LM is an (\mathcal{F}_t) -martingale. This is a consequence of Fubini's provided we can justify

$$\mathbb{E}(L_t M_t R_s) = \int_{\mathbb{R}} d\mu(u) \mathbb{E}[L_t(u) M_t R_s]. \quad (4.46)$$

For this we need to estimate

$$\int_{\mathbb{R}} d\mu(u) \mathbb{E}(|L_t(u) M_t R_s|). \quad (4.47)$$

By Cauchy-Schwarz the square of expression (4.47) is bounded by

$$\|R\|_{\infty}^2 \int_{\mathbb{R}} d\mu(u) \mathbb{E}(|L_t(u)|^2) \int_{\mathbb{R}} d\mu(u) \mathbb{E}(|M_t|^2) \leq |\mu|(\mathbb{R})^2 \|R\|_{\infty}^2 \mathbb{E}(|M_T|^2) \sup_{t \leq T; u \in \mathbb{R}} \mathbb{E}(|L_t(u)|^2)$$

(4.47) follows by (4.45). This finally shows that the expression (4.29) in the statement of Theorem 4.11 is an FS type decomposition which could be theoretically complex.

5. It remains to prove that the decomposition is real-valued. Let (H_0, ξ, L) and $(\overline{H}_0, \overline{\xi}, \overline{L})$ be two FS decomposition of H . Consequently, since H and (S_t) are real-valued, we have

$$0 = H - \overline{H} = (H_0 - \overline{H}_0) + \int_0^T (\xi_s - \overline{\xi}_s) dX_s + (L_T - \overline{L}_T),$$

which implies that $0 = \text{Im}(H_0) + \int_0^T \text{Im}(\xi_s) dX_s + \text{Im}(L_T)$. By Theorem 2.5, the uniqueness of the real-valued Föllmer-Schweizer decomposition yields that the processes $(H_t), (\xi_t)$ and (L_t) are real-valued.

□

5 The error in the quadratic minimization problem

Let $H \in \mathcal{L}^2$. The problem of minimization of the quadratic error given in Definition 2.1 is strongly connected with the FS decomposition. We evaluate now the error committed by the mean-variance hedging procedure. In the following lemma, we first calculate $\langle L(u), L(v) \rangle$ for any $u, v \in \mathbb{R}$.

Lemma 5.1. *We have*

$$\langle L(u), L(v) \rangle_t = \int_0^t \epsilon_{t,T}(u) \epsilon_{t,T}(v) e^{i \int_t^T (\delta_r(u) + \delta_r(v)) d\Psi'_r(0)} d\Gamma_s(u, v), \quad (5.1)$$

where $(V_t(u))$ is the exponential martingale defined by $V_t(u) = e^{iuX_t} \epsilon_{t,T}(u)$, as introduced in (4.17) and

$$\Gamma_t(u, v) = \nu_t(u, v) - \int_0^t \delta_s(u) \delta_s(v) d(-\Psi''_s(0)), \quad \text{with} \quad (5.2)$$

$$\nu_t(u, v) = \Psi_t(u + v) - \Psi_t(u) - \Psi_t(v) \quad (5.3)$$

Proof. We have

$$\begin{aligned} L_t(u) &= H_t(u) - H_0(u) - \int_0^t \xi_r(u) dX_r \\ &= V_t(u) e^{i \int_t^T \delta_s(u) d\Psi'_s(0)} - e^{\Psi_T(u) + i \int_0^T \delta_s(u) d\Psi'_s(0)} - \int_0^t \xi_r(u) dM_r - \int_0^t \xi_r(u) dA_r. \end{aligned}$$

Using integration by parts and the fact that $t \mapsto \int_0^t \delta_s(u) d\Psi'_s(0)$ is continuous (since $t \mapsto \Psi'_t(0)$ is),

$$L_t(u) = \mathcal{M}_t(u) + \mathcal{A}_t(u), \quad (5.4)$$

where

$$\begin{aligned} \mathcal{M}_t(u) &= \int_0^t e^{i \int_r^T \delta_s(u) d\Psi'_s(0)} dV_r(u) - \int_0^t \xi_r(u) dM_r - e^{\Psi_T(u) + i \int_0^T \delta_s(u) d\Psi'_s(0)} \\ \mathcal{A}_t(u) &= -i \int_0^t e^{i \int_r^T \delta_s(u) d\Psi'_s(0)} V_r(u) \delta_r(u) d\Psi'_r(0) + i \int_0^t \xi_r(u) d\Psi'_r(0). \end{aligned} \quad (5.5)$$

We observe that $\mathcal{A}_t(u)$ is predictable, $\mathcal{A}_0(u) = 0$, $L_0(u) = 0$. By uniqueness of the decomposition of an (\mathcal{F}_t) -special semimartingale, we obtain

$$L(u)_t = \mathcal{M}_t(u) = \int_0^t Z_s(u) d(-\Psi''_s(0)) \quad (5.6)$$

Since $\langle O(u), M \rangle = 0$ where $O(u)$ was defined in (4.22), it follows

$$\langle V(u), M \rangle_t = \int_0^t Z_s(u) d\langle M \rangle_s, \quad \text{where} \quad Z_t(u) = \delta_t(u) V_t(u). \quad (5.7)$$

We need at this point to express the predictable covariation $\langle V(u), V(v) \rangle$, $\forall u, v \in \mathbb{R}$. For this we decompose the product $V(u)V(v)$ to obtain

$$V_t(u)V_t(v) = V_t(u+v)R_t(u, v), \quad (5.8)$$

where

$$R_t(u, v) = \frac{\epsilon_{t,T}(u)\epsilon_{t,T}(v)}{\epsilon_{t,T}(u+v)} = \exp\{-(\nu_T(u, v) - \nu_t(u, v))\}.$$

Since $(R_t(u, v))_{t \in [0, T]}$ is continuous, integrating by parts we obtain

$$V_t(u)V_t(v) = \int_0^t R_s(u, v) dV_s(u+v) + \int_0^t V_s(u+v) R_s(u, v) d\nu_s(u, v).$$

Since $(\int_0^t R_s(u, v) dV_s(u + v))_t$ is an (\mathcal{F}_t) -local martingale, it follows that

$$\langle V(u), V(v) \rangle_t = \int_0^t V_s(u + v) R_s(u, v) (d\Psi_s(u + v) - d\Psi_s(u) - d\Psi_s(v)) . \quad (5.9)$$

We come back to the calculus of $\langle L(u), L(v) \rangle_t$; (5.5) and (5.6) give

$$\begin{aligned} \langle L(u), L(v) \rangle_t &= \left\langle L(u), \int_0^t e^{i \int_r^T \delta_s(v) d\Psi'_s(0)} dV_r(v) \right\rangle_t \\ &= \left\langle \int_0^t e^{i \int_r^T \delta_s(u) d\Psi'_s(0)} dV_r(u), \int_0^t e^{i \int_r^T \delta_s(v) d\Psi'_s(0)} dV_r(v) \right\rangle_t \\ &\quad - \left\langle \int_0^t \xi_r(u) dM_r, \int_0^t e^{i \int_r^T \delta_s(v) d\Psi'_s(0)} dV_r(v) \right\rangle_t \\ &= \int_0^t e^{i \int_r^T (\delta_s(u) + \delta_s(v)) d\Psi'_s(0)} d \langle V(u), V(v) \rangle_r - \int_0^t \xi_r(u) e^{i \int_r^T \delta_s(v) d\Psi'_s(0)} d \langle M, V(v) \rangle_r \end{aligned}$$

Using (5.7), (5.8) and (5.9), we obtain

$$\begin{aligned} \langle L(u), L(v) \rangle_t &= \int_0^t e^{i \int_r^T (\delta_s(u) + \delta_s(v)) d\Psi'_s(0)} V_r(u + v) R_r(u, v) d\nu_r(u, v) \\ &\quad - \int_0^t V_r(u) V_r(v) e^{i \int_r^T (\delta_s(u) + \delta_s(v)) d\Psi'_s(0)} \delta_r(u) \delta_r(v) d(-\Psi''_r(0)) \\ &= \int_0^t e^{i \int_r^T (\delta_s(u) + \delta_s(v)) d\Psi'_s(0)} \epsilon_{r,T}(u) \epsilon_{r,T}(v) \left[d\nu_r(u, v) + \delta_r(v) \delta_r(u) d\Psi''_r(0) \right] , \end{aligned}$$

which concludes the proof. \square

Now we can evaluate the error committed by the mean-variance hedging procedure described at Section 5.

Theorem 5.2. *Let $X = (X_t)_{t \in [0, T]}$ be a semimartingale with independent increments with log-characteristic function Ψ . Then the variance of the hedging error equals $J_0 := \int_{\mathbb{R}^2} d\mu(u) d\mu(v) J_0(u, v)$ where*

$$J_0(u, v) = \int_0^T \exp \left(\int_t^T \left(\frac{d\Psi'_s(0)}{d\Psi''_s(0)} \right)^2 d\Psi''_s(0) + i \int_t^T (\delta_s(u) + \delta_s(v)) d\Psi'_s(0) \right) \epsilon_{t,T}(u) \epsilon_{t,T}(v) d\Gamma_t(u, v) ,$$

where $\epsilon_{t,T}$ is defined in (4.1), δ_t is defined in (4.2) and Γ is defined in (5.2).

Proof. Theorem 2.6 implies that the variance of the hedging error equals

$$\mathbb{E} \left(\int_0^T \exp \{ -(K_T - K_s) \} d \langle L \rangle_s \right) , \quad (5.10)$$

where K was defined in Definition 2.3. By (3.10), it follows that $K_t = \int_0^t \left(\frac{d\Psi'_s(0)}{d\Psi''_s(0)} \right)^2 d(-\Psi''_s(0))$. We come back to expression (5.1) given in Lemma 5.1. It gives

$$\langle L(u), L(v) \rangle_t = C^1(u, v, t) + C^2(u, v, t) \quad (5.11)$$

where

$$C^1(u, v, t) := \int_0^t e^{i \int_s^T (\delta_r(u) + \delta_r(v)) d\Psi'_r(0)} \epsilon_{s,T}(u) \epsilon_{s,T}(v) e^{i(u+v)X_s} d\nu_s(u, v) \quad (5.12)$$

$$C^2(u, v, t) := \int_0^t e^{i \int_s^T (\delta_r(u) + \delta_r(v)) d\Psi_r'(0)} \epsilon_{s,T}(u) \epsilon_{s,T}(v) \delta_s(u) \delta_s(v) e^{i(u+v)X_s} d\Psi_s''(0). \quad (5.13)$$

We need to show that

$$\mathbb{E} \left(\int_{\mathbb{R}^2} d|\mu|(u) d|\mu|(v) \| \langle L(u), L(v) \rangle \|_{var} \right) < \infty. \quad (5.14)$$

Observe that

$$\|C^1(u, v, \cdot)\|_{var} = \int_0^T |H_s(u)H_s(v)| d|\nu_s|(u, v).$$

Let us consider first the term involving the measure ν . Notice that

$$d\nu_s(u, v) = \left[-uvc_s + \int_{\mathbb{R}} (e^{iux} - 1)(e^{ivx} - 1)F_s(dx) \right] d(-\Psi_s''(0)),$$

Then, we obtain the following upper bound

$$\begin{aligned} d|\nu_s|(u, v) &\leq \left[uv|c_s| + 2 \int_{\mathbb{R}} \left| \sin \frac{ux}{2} \right| \left| \sin \frac{vx}{2} \right| F_s(dx) \right] d(-\Psi_s''(0)) \\ &\leq \frac{1}{2} \left[(u^2 + v^2)c_s + 2 \int_{\mathbb{R}} \left[\left(\sin \frac{ux}{2} \right)^2 + \left(\sin \frac{vx}{2} \right)^2 \right] F_s(dx) \right] d(-\Psi_s''(0)) \\ &\leq \frac{1}{2} [\gamma_s(u) + \gamma_s(v)], \end{aligned}$$

where γ_s is defined in (4.26). we obtain by setting $\gamma_s(u, v) = (u^2 + v^2)\frac{c_s}{2} + \int_{\mathbb{R}} (\sin \frac{ux}{2})^2 (\sin \frac{vx}{2})^2 F_s(dx)$ Now, using inequality (4.33) yields

$$\begin{aligned} \|C^1(u, v, \cdot)\|_{var} &= \frac{1}{2} \exp(2K_T) \int_0^T [\gamma_s(u) + \gamma_s(v)] \exp \left\{ -\frac{1}{4} \int_s^T [\gamma_r(u) + \gamma_r(v)] d(-\Psi_r''(0)) \right\} d(-\Psi_s''(0)) \\ &= 2 \exp(2K_T) \left(1 - \exp \left\{ -\frac{1}{4} \int_0^T [\gamma_r(u) + \gamma_r(v)] d(-\Psi_r''(0)) \right\} \right) \\ &\leq 2 \exp(2K_T), \end{aligned}$$

which implies, by the fact that μ is finite

$$\int_{\mathbb{R}^2} d|\mu|(u) d|\mu|(v) \|C^1(u, v, \cdot)\|_{var} \leq 2 \exp(2K_T) |\mu|(\mathbb{R})^2. \quad (5.15)$$

Now, using (4.35),

$$\begin{aligned} \|C^2(u, v, \cdot)\|_{var} &\leq \int_0^T |\xi_s(u)| |\xi_s(v)| d(-\Psi_s''(0)) \\ &\leq 2 \exp(2K_T) \int_0^T \sqrt{\gamma_s(u)\gamma_s(v)} \exp \left\{ -\frac{1}{4} \int_s^T (\gamma_r(u) + \gamma_r(v)) d(-\Psi_r''(0)) \right\} d(-\Psi_s''(0)) \\ &\leq -4 \exp(2K_T) \int_0^T \left(-\frac{1}{4} \right) (\gamma_s(u) + \gamma_s(v)) \exp \left\{ -\frac{1}{4} \int_s^T (\gamma_r(u) + \gamma_r(v)) d(-\Psi_r''(0)) \right\} d(-\Psi_s''(0)) \\ &\leq 4 \exp(2K_T). \end{aligned} \quad (5.16)$$

Finally (5.16) implies

$$\int_{\mathbb{R}^2} \|C^2(u, v, \cdot)\|_{var} d|\mu|(u) d|\mu|(v) \leq 4 \exp(2K_T) |\mu|(\mathbb{R})^2. \quad (5.17)$$

Since K_T is deterministic, (5.17) and (5.15) imply (5.14). Previous considerations allow to prove that

$$\langle L \rangle_t = \int_{\mathbb{R}^2} d\mu(u) d\mu(v) \langle L(u), L(v) \rangle_t. \quad (5.18)$$

For this, it is enough to show that

$$L_t^2 - \int_{\mathbb{R}^2} d\mu(u)d\mu(v) \langle L(u), L(v) \rangle_t \quad (5.19)$$

produces an (\mathcal{F}_t) -martingale. First (5.14) shows that the second term in (5.19) is well-defined. By (4.45) and Fubini's, (5.19) gives $\int_{\mathbb{R}^2} d\mu(u)d\mu(v)(L_t(u)L_t(v) - \langle L(u), L(v) \rangle_t)$. By similar arguments as point 4.(a) in the proof of Theorem 4.11, using the fact that $(L_t(u))$ is a martingale and applying Fubini's, we are able to show that (5.19) defines a martingale. According to (5.10), the last step of the proof consists in evaluating the expectation of $\int_0^T \exp\{-(K_T - K_s)\} d\langle L \rangle_s$ taking into account (5.18) and Lemma 5.1. \square

6 Examples

6.1 The Gaussian examples

We refer here to the toy model introduced at Section 3.3.1. We suppose that $X_t = \gamma(t) + W_{\psi(t)}$ with ψ increasing, $d\gamma \ll d\psi$ and $\frac{d\gamma}{d\psi} \in \mathcal{L}^2(d\psi)$. This guarantees the (SC) property because of Proposition 3.7.2. Given f and μ expressed via (4.7), the FS-decomposition of $H = f(X_T)$ is provided by Theorem 4.11 with

$$\Psi'_t(0) = i\gamma(t), \quad \delta_t(u) = iu, \quad \text{and} \quad \epsilon_{t,T}(u) = \exp[iu(\gamma(T) - \gamma(t)) - \frac{u^2}{2}(\psi(T) - \psi(t))],$$

which yields

$$H_t(u) = \exp[2iu(\gamma(T) - \gamma(t)) - \frac{u^2}{2}(\psi(T) - \psi(t))]e^{iuX_t} \quad \text{and} \quad \xi_t(u) = iuH_t(u)$$

According to Lemma 5.1, we can easily show that $\Gamma_t(u, v) \equiv 0$, for all $t \in [0, T]$, $u, v \in \mathbb{R}$. Consequently, the variance of the hedging error is zero.

6.2 The Lévy case

Let X be a square integrable Lévy process, with characteristic function $\Psi_t(u)$ where $\Psi_t(u) = t\Psi(u)$. It is always a semimartingale since $\Psi \rightarrow e^{i\Psi_t(u)}$ has bounded variation, see Theorem 4.14 of [17]. By Remark 3.4, Ψ is of class $C^2(\mathbb{R})$. We suppose that $\Psi''(0) \neq 0$. We have

$$\frac{d\Psi'_t(0)}{d\Psi''_t(0)} = \frac{\Psi'(0)}{\Psi''(0)}.$$

Condition (SC) is verified taking into account Proposition 3.7. In conclusion, we can apply Theorem 4.11 taking into account (4.1) and (4.2), we obtain $V_t(u) = \exp((T-t)\Psi(u))e^{iuX_t}$,

$$H_t(u) = \exp\left((T-t)\left(\Psi(u) + \frac{\Psi'(u) - \Psi'(0)}{\Psi''(0)}\Psi'(0)\right)\right)e^{iuX_t} \quad \text{and} \quad \xi_t(u) = H_t(u)i\frac{\Psi'(u) - \Psi'(0)}{\Psi''(0)}.$$

The factor $\Gamma_t(u, v)$ appearing in Lemma 5.1 gives $\Gamma_t(u, v) = t\Gamma(u, v)$ and

$$\Gamma(u, v) = (\Psi(u+v) - \Psi(u) - \Psi(v)) + \frac{(\Psi'(v) - \Psi'(0))(\Psi'(u) - \Psi'(0))}{(-\Psi''(0))}.$$

In particular, when X is a Poisson process we have $\Gamma(u, v) \equiv 0$. This shows, as expected, that $\langle L(u), L(v) \rangle = 0, \forall u, v \in \mathbb{R}$.

6.3 Wiener integral of Lévy processes

With the same notations as in subsection 3.3.3, we consider a square integrable Lévy process $\Lambda = (\Lambda_t)_{t \in [0, T]}$ such that $\Lambda_0 = 0$ and $\text{Var}(\Lambda_1) \neq 0$. Let $\gamma : [0, T] \rightarrow \mathbb{R}$ be a bounded Borel function. We set $X_t = \int_0^t \gamma_s d\Lambda_s$, $t \in [0, T]$. For $u \in \mathbb{R}$, $t \in [0, T]$, we have the following quantities. According to the observations below Remark 3.13

$$\epsilon_{t,T}(u) = \exp \left(\int_t^T \Psi_\Lambda(u\gamma_s) ds \right) \quad (6.20)$$

$$\delta_{t,T}(u) = -i \frac{\Psi'_\Lambda(u\gamma_t) - \Psi'_\Lambda(0)}{\Psi''_\Lambda(0)}. \quad (6.21)$$

Remark 6.1. If $\Lambda = W$ then there is a Brownian motion \tilde{W} such that $X_t = \tilde{W}_{\Psi(t)}$ with $\Psi(t) = \int_0^t \gamma_s^2 ds$. This was the object of Section 6.1.

6.4 Representation of some contingent claims by Fourier transforms

In general, it is not possible to find a Fourier representation, of the form (4.7), for a given payoff function which is not necessarily bounded or integrable. Hence, it can be more convenient to use the bilateral Laplace transform that allows an extended domain of definition including non integrable functions. We refer to [8], [25] and more recently [11] for such characterizations of payoff functions. It should be certainly possible to extend our approach replacing the Fourier transform with the bilateral Laplace transform. However, to illustrate the present approach restricted to payoff functions represented as classical Fourier transforms, we give here one simple example of such representation extracted from [11]. The payoff of a *self quanto put option* with strike K is

$$f(x) = e^x(K - e^x)_+ \quad \text{and} \quad \hat{f}(u) = \int_{\mathbb{R}} e^{iux} f(x) dx = \frac{K^{2+iu}}{(1+iu)(2+iu)}.$$

In this case μ admits a density which is proportional to \hat{f} which is integrable.

Appendix

Proof of Proposition 3.10. In the sequel, we will make use of Lemma 3.12 of [13] in a fairly extended generality.

Lemma 6.2. Let \mathcal{N} be a complete metric space and μ and ν are two non-negative Radon non-atomic measures. We suppose the following:

1. $\mu \ll \nu$;
2. $\mu(I) \neq 0$ for every open ball I of \mathcal{N} .

Then $h := \frac{d\mu}{d\nu} \neq 0$ ν a.e. In particular μ and ν are equivalent.

1. If there are no deterministic increments then setting $d\mu(t) = -d\Psi_t''(0)$ and $d\nu(t) = da_t$, it follows that da_t is equivalent to $-d\Psi_t''(0)$, because of Lemma 6.2. Consequently the result is established.

2. Suppose that not all the increments are non-deterministic. We decompose $E := [0, T] = E_R \cup E_R^C$ where

$$E_R = \{t \in [0, T] | \text{Var}(X_{(t+\varepsilon)\wedge T} - X_t) > 0, \forall \varepsilon > 0 \text{ or } \text{Var}(X_t - X_{(t-\varepsilon)_+}) \neq 0 \forall \varepsilon > 0\},$$

and its complementary

$$E_R^C = \{t \in [0, T] | \exists \varepsilon > 0, \text{Var}(X_{(t+\varepsilon)\wedge T} - X_{(t-\varepsilon)_+}) = 0\}.$$

Without restriction to generality we can suppose that $T \in E_R$. Since E_R^C is an open subset of $[0, T]$, it can be decomposed into a union $\bigcup_{n \in \mathbb{N}} I_n$ of open (disjoint) intervals of $[0, T]$. We denote $a_n = \inf I_n, b_n = \sup I_n$. Clearly a_n and b_n belong to E_R and to its boundary, since E_R is closed. We define on E a semidistance d such that $d(u, v) = \text{Var}(X_v - X_u)$. The equivalence relation \mathcal{R} on E defined setting $x \mathcal{R} y$ if and only if $d(x, y) = 0$, produces the following equivalence classes:

$$\{t\}, t \in \text{int} E_R, \quad I_n, n \in \mathbb{N}.$$

The quotient E/d can be identified with family of *typical representatives* $E_d = \text{int} E_R \cup \{b_n, n \in \mathbb{N}\}$.

We denote by $\tilde{a}_t = \int_{[0, t] \cap E_R} da_s$. The proof of Proposition 3.10 will be concluded if the two lemmas below hold.

Lemma 6.3. (a) $\int_{E_R^C} d(-\Psi_t''(0)) = 0$.

(b) $\Psi_t(u)$ (resp. $\Psi_t''(u)$) is absolutely continuous with respect to $d\tilde{a}$, for every $u \in \mathbb{R}$.

(c) $\int_{E_R^C} d\tilde{a}_t = 0$.

Lemma 6.4. $d\tilde{a}$ is equivalent to $d(-\Psi_t''(0))$.

Proof of Lemma 6.3. (a) Since each I_n is precompact, it can be recovered by a countable sequence of subintervals of the type $]t_n - \varepsilon_n, t_n + \varepsilon_n[$. So, by definition of E_R^C , we have $\int_{I_n} d(-\Psi_t'')(0) = 0$. The item follows then because E_R^C is the union of countable intervals.

(b) It is enough to show that for every B Borel subset of E_R^C we have

$$\int_B \xi_s(u) da_s = \int_B \eta_s(u) da_s = 0,$$

where $\eta_s(u)$ (resp. $\xi_s(u)$) was introduced in (3.2) (resp. (3.8)). We only treat the $\Psi_t(u)$ case, the other one being similar. By Proposition 3.9, if $X_b - X_a$ is deterministic then X and in particular $t \mapsto \Psi_t(u)$ is constant on $[a, b]$. Consequently

$$\left| \int_{E_R^C} \eta_s(u) da_s \right| \leq \sum_{n \in \mathbb{N}} \left| \int_{I_n} \eta_s(u) da_s \right| = \sum_{n \in \mathbb{N}} \left| \int_{I_n} d\Psi_t(0) \right| = 0.$$

(c) It is a consequence of the definition of \tilde{a} .

□

Proof of Lemma 6.4. $\mathcal{N} := E_d$ is a complete metric space equipped with the distance, inherited from E , still denoted by d . We define $d\mu$ (resp. $d\nu$) the measure on the Borel σ -algebra of \mathcal{N} obtained by restriction from $-d\Psi_t''(0)$ (resp. $d\tilde{a}_t$). This is possible since by items (a) and (c) of Lemma 6.3

$$\int_{E_R^C} d(-\Psi_t''(0)) = \int_{E_R^C} d\tilde{a}_s = 0.$$

Condition 1 of Lemma 6.2 is verified by item (b). Concerning Condition 2. of the same lemma, let $t_0 \in E_d \subset E_R$ and $B(t_0)$ an open ball centered at t_0 . Obviously $\mu(B(t_0)) > 0$. By Lemma 6.2, $\nu \sim \mu$ on the Borel σ -algebra of E_d and the result follows. \square

This concludes the proof of Proposition 3.10. \square

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