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Remote triggering of air-gap discharge by a femtosecond laser filament and postfilament at distances up to 80 m

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Abstract. We experimentally observed a laser-induced remote high-voltage discharge triggering between two needle electrodes with half-a-cm spacing. The discharge was initiated by a 744-nm, 90-fs, 6-mJ laser pulse undergoing filamentation in air. For the direct voltage below the self-breakdown threshold, triggering of air-gap discharge was synchronized with 10-Hz laser repetition rate and occurred between 40 and 80 m of the propagation path. No discharge guiding was observed. The experimentally registered and simulated remote triggering probability was above 80% in the range of 40–65 m from laser output, and about 50% in the range of 65–80 m. The probability decreases as the postfilament hot spot diverges with simultaneous increase of stochastic laser beam wandering.

Laser induced discharge has a lot of applications in different areas of science and technology. Discharge initiation is crucial in the systems where precise synchronization of different electrical processes is needed. For instance, laser-triggered switches are very useful in laser systems where the active medium is pumped by independent sources (as in KrF laser facility GARPUN¹). Indeed, the spark gap triggered by a femtosecond laser filament in air combines the advantages of large interelectrode air gaps capable of switching high voltages and low subnanosecond jitter of narrow spark gaps². Likewise laser induced discharges seem to be very promising for the fuel ignition of rocket engines³ and cryogenic thrusters^{4,5}. Another important feature of laser induced electric discharges is their ability to channel a lot of electrical energy at high currents⁶. This effect can be applied to solve the problem of friction in high-speed train's power supply⁷: in Ref. 8 such discharges were suggested to replace the pantographs of the trains. Laser induced plasma not only commutes the electrodes by ignition of the electric discharge, but also gives an opportunity of controlling the time interval of vacuum gap closure⁹.

After the discovery of self-focusing and laser air breakdown, the laser pulses were proposed to trigger and guide electric discharges^{10–14}. The use of laser pulses for lightning protection¹⁵ was patented¹⁶ and carried out with a combination of two-arm CO₂, Nd:glass, and Nd:YAG high power lasers with tens of nanosecond pulse duration¹⁷. The CO₂ laser induced the plasma on a dielectric target at the top of the lightning tower, while the fourth harmonic of the YAG laser

produced weakly ionized plasma channels far above the target to guide a leader towards thunderclouds¹⁷. Only two cases of laser-triggered lightning were recorded from the extensive material accumulated¹⁷ in the years 1997–99.

The development of the Chirped Pulse Amplification technology led to the observation of filamentation¹⁸ of collimated femtosecond laser pulses in 1995. In the same year, new schemes based on the extended plasma channels formed by femtosecond laser pulses were proposed for lightning protection^{19,20}. Several groups worldwide performed experiments to trigger high-voltage discharges with femtosecond pulses^{21–26}. Laboratory experiments evidenced triggered and guided discharges over meter-scale distances^{21,23,27,28}. Experiments on lightning in the atmosphere rely on the potentially significant ionization induced by multiple filaments over a typical length of 100 m, a few hundreds of meters above the ground. A statistically significant number of electric events synchronized with the laser pulses of high-power *Teramobile*²⁹ facility was reported after real-scale experiments³⁰. This research is continued nowadays, and a laser lightning rod based on the plasma and the low-density channel produced by a filament of 1.03- μm , ~ 1 -ps, Joule-level and high repetition rate laser is under investigation³¹.

An overwhelming majority of research on discharge triggering was implemented with focused laser beams, which limited the range of application of such switches. The plasma formed in the focus either connects the electrodes with the conducting plasma channel^{21,25,26} or heats the gas, leaving a low-density air channel after expansion of the hot region^{22,32}. Both cases limit the range of possible discharge triggering to the region with a sufficient number of free electrons. This is very convenient in the case of focusing geometry, because external geometrical focusing determines the position of a fil-

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ament and, consequently, that of the plasma channel. In the collimated geometry, beam pointing instability as well as air turbulence lead to fluctuations of the filament start³³ and make it hard to obtain a plasma channel at the desired distance.

At the same time, comparatively weak intensities can trigger discharge without guiding it¹⁹. Both triggering and guiding can occur for a lower voltage than the self-breakdown voltage. Most often, a discharge can be simply triggered by generating a localized plasma in the gap between the two electrodes, for instance, by focusing a femtosecond laser pulse²². In contrast, a laser guided discharge requires a predefined path connecting the electrodes, where the mean free path of electrons is larger than in the surrounding region. The hot column of air left by a femtosecond pulse in its wake plays this role. In this case, a guided discharge is observed for a shorter delay (a few tens of nanoseconds, depending on the gap and applied voltage^{22,34}) after the triggering pulse than for an unguided discharge (up to several microseconds²²). So, depending on the filament geometry relative to the electrodes, the discharge may be guided^{19,21–23,35}, even along a curve^{32,36}, partly guided^{23,27}, or just triggered²².

In this work we remotely *trigger* an electric discharge between two needle electrodes placed along a 100-m-long path with a femtosecond laser pulse propagating in a single filament regime. The discharge was observed between electrodes located anywhere within the high-intensity region, a distance from 40 m to 80 m away from the laser. At the distances 45–60 m, the probability of the discharge is 80%, for longer distances it decreases while remaining about 50%. The orientation of the needle electrodes relative to the beam propagation direction did not affect the discharge occurrence. The discharge triggering at a certain distance is associated with a high-intensity laser beam area of about 1 mm^2 [Fig. 1(a)] hitting the cathode and generating photoelectrons from its metal surface as predicted in Ref. 37. Three-dimensional carrier wave resolved numerical simulations on a 100-m path in air show that the intensity threshold of $\sim 0.5 \text{ TW/cm}^2$ should be overcome to trigger a discharge in the long-range experiment.

The experiments were carried out using the pulses from a commercial Ti:Sapphire laser system (Avesta Ltd.) with a central wavelength of 744 nm, a pulse duration of 90 fs (FWHM), pulse energy up to 6 mJ and a repetition rate of 10 Hz. The beam was elliptical with the ratio of 0.4 between the semi-axes and the length of the major axis ($\text{FW}e^{-1}\text{M}$) of 8 mm. This beam was directed to the $\sim 100\text{-m}$ corridor inside the building using several highly reflective interference mirrors (the details are given in Ref. 38). The mirror system introduced a slight divergence to the laser beam (the focal length of the system was estimated as -50 m) thus increasing the distance of the beam collapse to 40–50 m. According to our previous analysis³⁸, the single filament consisted in a short (about 0.3-m) region of high intensity located at $z \sim 50 \text{ m}$ with plasma of density in excess of 10^{16} cm^{-3} , followed by a long postfilament channel without plasma but an intensity in excess of 0.5 TW/cm^2 . Single “self-cleaned” hot spot surrounded by wide elliptical background was observed from 50 to 100 m, see Ref. 38 and Fig. 1(a).

The discharge triggering system consisted of two needle

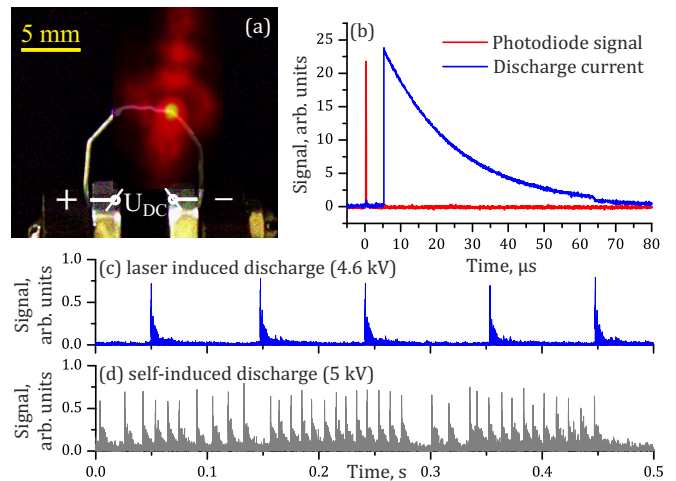


FIG. 1. (a) Photo of the needle electrodes connected by a discharge. The profile of beam luminescence from a paper screen in false color scale taken by CCD-camera at 75 m from the laser system is superimposed over the cathode in the same spatial scale. The postfilament hot spot (a few TW/cm^2) appears as a 1.2-mm yellow circle. (b) Photodiode signal (red curve) and discharge current (blue) registered by oscilloscope. (c,d) Audio signal recorded by microphone while multiple successive (c) laser- and (d) self-induced discharges occurred.

steel electrodes whose tips are facing each other, separated by 4–5 mm [Fig. 1(a)]. The diameter of the electrode’s needles was about 0.5 mm. The surface of the tips is irregular and sharp points can appear and disappear due to the damage caused by high-voltage discharge. The laser pulse propagated either perpendicular [Fig. 1(a)] or parallel to the line connecting the tips of the electrodes and hit the cathode.

The high-voltage supply supported the direct voltage up to 10 kV, which dropped down during the discharge occurrence only and regained after the discharge current termination. When the voltage between the electrodes was higher than the self-breakdown one of $\sim 4.8 \text{ kV}$, the arc appeared independently of a laser shot, that is, even if we blocked the beam. The self-induced discharge events were generated with the frequency of $(80 \pm 30) \text{ Hz}$ [Fig. 1(d)], which is much larger than the laser repetition rate of 10 Hz. To study the laser-induced discharge, we set the voltage 100–500 V below the self-breakdown threshold. In that case, the discharge events appeared at the laser repetition rate [Fig. 1(c)].

The discharge current was detected by an oscilloscope. Behind the electrodes we placed a paper screen to block the laser beam while a photodiode responded to the light scattered from this screen. The signal from the photodiode was a synchronization pulse for the oscilloscope. We considered the electric discharge as the laser driven one if only the discharge current was synchronized with the signal from the photodiode [see Fig. 1(b), a typical delay between the photodiode signal and the current onset was less than $25 \mu\text{s}$] and checked regularly the absence of discharge when we blocked the laser beam.

Discharge triggering was observed at the distances z from 40 to 80 m from the laser system output [Fig. 2(a), dashed rectangle]. The measured beam radius at the specific distances

of the discharge observation is marked with stars in Fig. 2(a). The discharge occurred when the laser beam hit the cathode for any orientation of the line connecting the needle tips, either perpendicular or parallel to the laser beam propagation direction. When the line connecting the tips was parallel to the propagation direction and the anode was closer to the laser (so, that it shadowed the cathode), the discharge was triggered less often (approximately once in 10 shots) than in all the other configurations when the laser beam hit the cathode directly. The straightening of the arc, which is the characteristic indication of the guided discharge^{19,35}, was never observed. Thus, at all distances z studied, the discharge was triggered but not guided, see Fig. 1(a).

In the range of distances $z = 45\text{--}60$ m the discharge was observed at almost every laser shot for an applied voltage close to the breakdown voltage of ~ 4.8 kV. At $z = 54$ m, we measured, out of 50 laser shots, the probability to initiate the discharge as a function of the applied voltage [Fig. 3(a)]. For the voltages of 4.6–4.7 kV this probability reached 80% revealing reasonable efficiency for the laser-triggered remote discharge.

Numerical simulations in the conditions of our experiment on a 100-m path were carried out to quantify the pulse intensity distribution in the filament and postfilamentation channel. We used the Forward Maxwell equation (FME, paraxial propagation equation for carrier-wave-resolved fields^{39,40}) in axially-symmetrical geometry $(t, r) + z$ due to the moderate pulse powers in our experiment. The nonlinear source term in FME includes transient photocurrent as well as the third-order instantaneous and delayed nonlinearities. The details about numerical grids required to capture finely the supercontinuum that diverges in the transverse domain and get delayed in time can be found in Ref. 38.

These simulations in the conditions of our long-range experiment demonstrate that in the 40-m discharge triggering zone the pulse peak intensity is higher than 0.5 TW/cm^2 . This numerically obtained intensity range agrees well with the intensity estimated from the experiment [Fig. 2(b), cf. the black solid curve to dots and stars]. We substituted the experimentally measured beam radius r [Fig. 2(a), stars and dots] into the equation $I = 2W\sqrt{\ln 2}/(\pi^{3/2}r^2\tau)$ and for the pulse energy of $W = 6$ mJ as well as the initial FWHM pulse duration $\tau = 90$ fs obtained the intensity I shown by dots and stars in [Fig. 2(b)]. This value is by far lower than the clamped intensity⁴¹ of $\sim 50 \text{ TW/cm}^2$ and is not enough to ionize air efficiently. In the simulations, the discharge triggering zone, except for the 0.3-m filament at $z \approx 48$ m, is characterized by the intensity in the range from 0.5 to 10 TW/cm^2 . The corresponding range of plasma density is from $\sim 10^3$ to $\sim 10^{11} \text{ cm}^{-3}$, respectively (cf. with a typical natural ion density of $10^2\text{--}10^4 \text{ cm}^{-3}$ in air in urban areas⁴²). Therefore, except for the 0.3-m plasma region, the laser pulse does not generate enough electrons to heat the air and does not create a straight preferred path for the discharge leader by the conducting plasma channel²³. Heating by the discharge arc itself does not happen since the post-arc temperature decay is faster^{43,44} than 100 ms.

The central high-intensity part of the laser beam propagating in the discharge triggering zone is 1–2 mm in diameter and

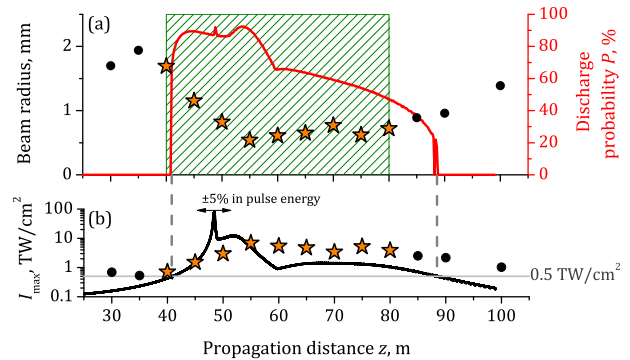


FIG. 2. (a) Dots and stars represent the measured dependence of the beam radius along the propagation direction z (left axis). Stars represent the positions where the discharge was reliably triggered. In accordance with confident triggering events, the dashed rectangle forms a 40-m discharge triggering zone. Red solid curve (right axis) is the calculated probability of the discharge triggering. (b) Dots and stars represent the intensity dependence on the distance z estimated from the measured beam radius in panel (a). Solid curve is the intensity obtained from the 3D+time simulations on a 100-m propagation path. The initial pulse energy fluctuations in the experiment introduce a ± 3 m displacement of this curve along the z direction.

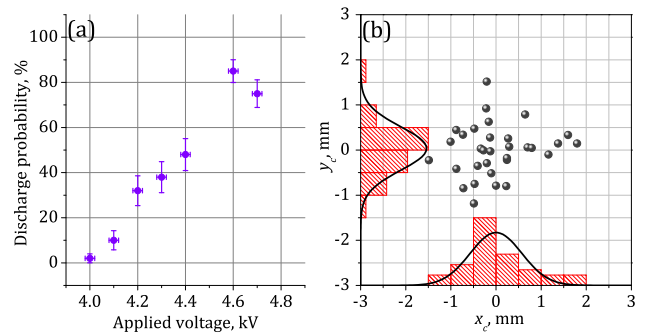


FIG. 3. (a) Experimentally measured dependence of discharge probability on the applied voltage at the distance $z = 54$ m. (b) Measured wandering of the beam center $\{x_c, y_c\}$ at the distance $z = 60$ m (spheres). The number of laser shots, for which the beam center $\{x_c, y_c\}$ falls within a 0.5-mm range along either x or y -axis (red bars, out of scale), and Gaussian fits of these distributions (black curves).

preserves this transverse size from 50 m to 80 m [Fig. 2(a), stars]. The physical reason for this notably stable hot spot diameter is the nonlinear transformation of a femtosecond pulse in the postfilament^{38,45–48}. The peak power of a femtosecond pulse in the postfilament is about the critical power for self-focusing in air. The laser beam in the postfilament does not diverge over a long distance, preserves high intensity and exceeds the 0.5-mm diameter of the needle electrode tip [Fig. 1(a)] by a factor of 2–4.

Therefore, the probability of hitting the cathode and initiating the discharge depends on the beam wandering in the transverse direction due to atmospheric turbulence, thermal lensing in amplification system etc. For quantitative characterization of the laser beam wandering, we recorded the beam

fluence distribution at $z = 60$ m in 32 successive laser shots. The “mass centers” $\{x_c, y_c\}$ of the measured fluence distributions [spheres in Fig. 3(b)] have the standard deviations along x and y axes of $\sigma_x = 1.15$ mm and $\sigma_y = 0.97$ mm, respectively.

The simulated spatio-temporal intensity distribution $I(t, r, z)$ on a 100-m path enables us to find the hot spot size, i.e., the largest radius r at each propagation position z for which the pulse intensity $I(r, t, z)$ exceeds the threshold value $I_{th} = 0.5$ TW/cm² chosen for the reason that the peak pulse intensity is above this value throughout the experimentally found discharge triggering zone, see Fig. 2.

We estimate the probability for the hot spot to illuminate the electrode tip and trigger the discharge by extracting from the simulated beam intensity distribution the area, in which the intensity exceeds I_{th} , and taking into account the beam wandering registered in the experiment [Fig. 3(b), spheres and histograms]. The hot spot center is assumed to be wandering according to a Gaussian distribution with the width increasing linearly with the propagation distance z : $\sigma(z) = \frac{1}{2}(\sigma_x + \sigma_y) \times z/60$ m ≈ 1 mm $\times z/60$ m. For the electrode tip with the diameter of 0.5 mm the resulting probability is presented in Fig. 2(a) by the red curve.

In agreement with the experiment, the calculated discharge probability P exceeds 80% in the 15-m-long region of the highest intensity of the light channel, which spans from 42 to 57 m [Fig. 2(a), the red curve]. Further on, for another 15 meters, up to the distance $z \approx 72$ m, the discharge is initiated less frequently, however, its probability is still higher than 50%. Beyond $z \approx 72$ m and up to $z \approx 90$ m the discharge probability decreases further due to the transverse beam wandering but remains above 20%. The effective discharge triggering radius, i.e. the largest radius of the hot spot, in which the intensity is above the threshold value $I_{th} = 0.5$ TW/cm², is 0.65 mm. Thus, the large beam background [Fig. 1(a), red color] does not contribute to the discharge triggering.

The nonlinear photoemission^{37,49–51}, including the one induced by a femtosecond pump^{52,53}, was observed at the intensity in the range of 1–10 GW/cm², i. e. 2–3 orders of magnitude lower than the threshold intensity $I_{th} = 0.5$ TW/cm² we used for calculating the hot spot radius and quantifying the probability of cathode tip hitting. As extracted from the simulations, the fluence of the laser pulse imposed on the cathode in our experiment is a factor of 2–10 below the material ablation threshold, except for the ~ 0.3 -m zone of high intensity [Fig. 2(b), the black curve at $z \approx 48$ m]. Indeed, the ablation can be observed at the intensity⁵⁴ close to $I_{th} = 0.5$ TW/cm², however, it requires a light fluence higher than⁵⁵ ~ 0.5 J/cm².

In conclusion, we show that the 90-fs, 744-nm, 6-mJ laser pulses at 10-Hz repetition rate can be used for triggering discharges at the positions of 40–80 m from laser output. The triggering efficiency approaches 80% in the ~ 15 -m-long region including the ~ 0.3 -m-long plasma filament and postfilament where the pulse intensity is about 5 TW/cm². Further down the propagation path the probability of the discharge triggering decreases due to the laser beam wandering.

Numerical simulations on a 100-m propagation path in air were used to quantify the pulse intensity in the discharge triggering zone. The threshold intensity of ~ 0.5 TW/cm² was

used to define the transverse size of the hot spot hitting the cathode tip. The joint usage of the ~ 0.5 TW/cm² threshold and the experimentally obtained laser beam center fluctuations yielded the values of the discharge probability, which agree with the experiment.

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DATA AVAILABILITY

The data that support the findings of this study are available from the corresponding author upon reasonable request.

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